

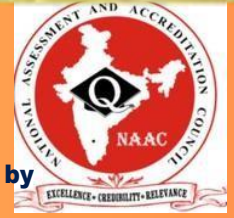


VSM COLLEGE OF ENGINEERING AUTONOMOUS

Accredited by NAAC with 'A' Grade - 3.23/4.00 CGPA

(Approved by AICTE, New Delhi and Permanently affiliated to JNTUK, Kakinada)

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The Ramchandrapuram Education Society (Estd. 1965)



Department of ELECTRICAL & ELECTRONICS ENGINEERING

ELECTROMAGNETIC FIELD THEORY

SUBJECT MATERIAL

YEAR : II SEMESTER : I

Regulation: VR23

Subject Code: VR2321201

Prepared by

Mrs. P. MANJUSHA

Associate Professor

EEE Department



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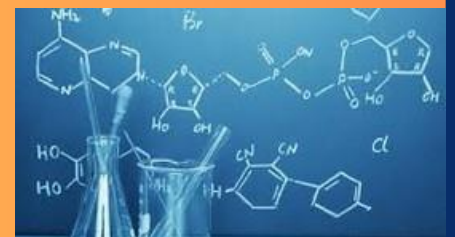
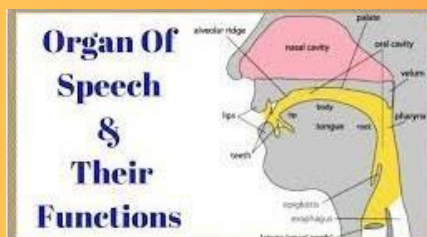


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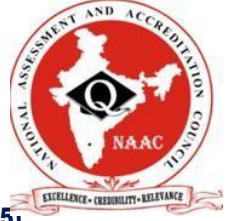
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Department of
ELECTRICAL & ELECTRONICS ENGINEERING
Subject Material
ELECTROMAGNETIC FIELD THEORY

II B.TECH I SEM

Regulation: VR23
Subject Code: VR2321201



VSM COLLEGE OF ENGINEERING
Ramachandrapuram-533255



VSM COLLEGE OF ENGINEERING [3B]

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II B.TECH I SEM

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SUB CODE: VR2321201

Electromagnetic Field Theory

Course Objectives:

- To study the production of electric field and potentials due to different configurations of static charges.
- To study the properties of conductors and dielectrics, calculate the capacitance of different configurations. Understand the concept of conduction and convection current densities.
- To study the magnetic fields produced by currents in different configurations, application of Ampere's law and the Maxwell's second and third equations.
- To study the magnetic force and torque through Lorentz force equation in magnetic field environment like conductors and other current loops.
- To develop the concept of self and mutual inductances and the energy stored.
- To study time varying and Maxwell's equations in different forms and Maxwell's fourth equation for the induced EMF.

Course Outcomes:

At the end of the course, student will be able to,

- CO1:** Compute electric fields and potentials using Gauss law/ solve Laplace's or Poisson's equations for various electric charge distributions.
- CO2:** Analyse the behaviour of conductors in electric fields, electric dipole and the capacitance and energy stored in dielectrics.
- CO3:** Calculate the magnetic field intensity due to current carrying conductor and understanding the application of Ampere's law, Maxwell's second and third law.
- CO4:** Estimate self and mutual inductances and the energy stored in the magnetic field.
- CO5:** Understand the concepts of Faraday's laws, Displacement current, Poynting theorem and Poynting vector.

UNIT - I

Vector Analysis:

Vector Algebra: Scalars and Vectors, Unit vector, Vector addition and subtraction, Position and distance vectors, Vector multiplication, Components of Coordinate Systems: Rectangular, Cylindrical and Spherical coordinate systems. Vector Calculus: Differential length, Area and Volume. Del operator, Gradient of a scalar, Divergence of a vector and Divergence theorem (definition only). Curl of a vector and Stoke's theorem (definition only), Laplacian of a scalar

Electrostatics:

Coulomb's law and Electric field intensity (E_f)- E_f due to Continuous charge distributions (line and surface charge), Electric flux density, Gauss's law (Maxwell's first equation, $\nabla \cdot \vec{D} = \rho_v$), Applications of Gauss's law, Electric Potential, Work done in moving a point charge in an electrostatic field (second Maxwell's equation for static electric fields, $\nabla \times \vec{E} = 0$), Potential gradient, Laplace's and Poisson's equations.

UNIT - II

Conductors – Dielectrics and Capacitance:

Behaviour of conductor in Electric field, Electric dipole and dipole moment Potential and EFI due to an electric dipole, Torque on an Electric dipole placed in an electric field, Current density-conduction and convection current densities, Ohm's law in point form, Behaviour of conductors in an electric field, Polarization, dielectric constant and strength, Continuity equation and relaxation time, Boundary conditions between conductor to dielectric, dielectric to dielectric and conductor to free space, Capacitance of parallel plate, coaxial and spherical capacitors, Energy stored and density in a static electric field.

UNIT - III

Magneto statics, Ampere's Law and Force in magnetic fields:

Biot-Savart's law and its applications viz. Straight current carrying filament, circular, square, rectangle and solenoid current carrying wire Magnetic flux density and Maxwell's second Equation ($\nabla \cdot \vec{B} = 0$), Ampere's circuital law and its applications viz. MFI due to an infinite sheet, long filament, solenoid, toroidal current carrying conductor, point form of Ampere's circuital law, Maxwell's third equation ($\nabla \times \vec{H} = \vec{J}$).

Magnetic force, moving charges in a magnetic field Lorentz force equation, force on a current element in a magnetic field, force on a straight and a long current carrying conductor in a magnetic field, force between two straight long and parallel current carrying conductors, Magnetic dipole, Magnetic torque, and moment.

UNIT - IV

Self and mutual inductance:

Self and mutual inductance – determination of self-inductance of a solenoid, toroid, coaxial cable and mutual inductance between a straight long wire and a square loop wire in the same plane – Energy stored and energy density in a magnetic field.

UNIT - V

Time Varying Fields:

Faraday's laws of electromagnetic induction, Maxwell's fourth equation ($\nabla \times \vec{E} = -\frac{d\vec{B}}{dt}$), integral and point forms of Maxwell's equations, statically and dynamically induced EMF, Displacement current, Modification of Maxwell's equations for time varying fields, Poynting theorem and Poynting vector.

Textbooks:

1. "Elements of Electromagnetics" by Matthew N O Sadiku, Oxford Publications, 7th edition, 2018.
2. "Engineering Electromagnetics" by William H. Hayt & John. A. Buck Mc. Graw-Hill 7th Edition. 2006.

Reference Books:

1. "Introduction to Electro Dynamics" by D J Griffiths, Prentice- Hall of India Pvt. Ltd, 2nd edition.
2. "Electromagnetic Field Theory" by Yaduvir Singh, Pearson India, 1st edition, 2011.
3. "Fundamentals of Engineering Electromagnetics" by Sunil Bhooshan, Oxford University Press, 2012.
4. Schaum's Outline of Electromagnetics by Joseph A. Edminister, Mahamood Navi 4th, Edition, 2014.

Online Learning Resources:

1. <https://archive.nptel.ac.in/courses/108/106/108106073/>
2. <https://nptel.ac.in/courses/117103065>

UNIT - I

Vector:-

$$\vec{A} = A_x \hat{a}_x + A_y \hat{a}_y + A_z \hat{a}_z$$

where $\hat{a}_x, \hat{a}_y, \hat{a}_z$ are unit vectors

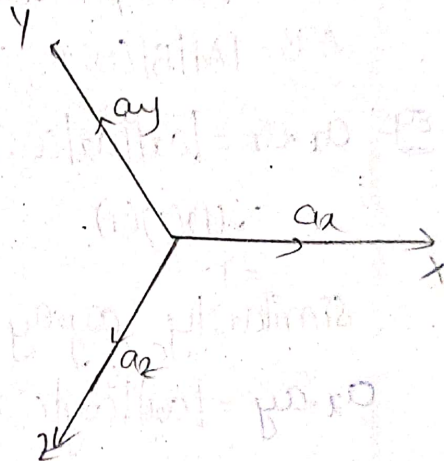
i.e. A_x, A_y, A_z are the co-efficients of a vector 'A' in x, y, z

$$|\vec{A}| = \sqrt{A_x^2 + A_y^2 + A_z^2}$$

i.e. unit vector \hat{a} is

$$\hat{a}_A = \frac{\vec{A}}{|\vec{A}|}$$

$$\hat{a}_A = \frac{A_x \hat{a}_x + A_y \hat{a}_y + A_z \hat{a}_z}{\sqrt{A_x^2 + A_y^2 + A_z^2}}$$



Eg:- $\vec{A} = 2\hat{a}_x + 3\hat{a}_y + 4\hat{a}_z$; $\vec{B} = 5\hat{a}_x + 4\hat{a}_y + 6\hat{a}_z$ Find distance between two vector.

Distance between two vector

$$\begin{aligned} \vec{R}_{AB} &= \vec{B} - \vec{A} \\ &= 5\hat{a}_x + 4\hat{a}_y + 6\hat{a}_z - 2\hat{a}_x - 3\hat{a}_y - 4\hat{a}_z \\ &= 3\hat{a}_x + \hat{a}_y + 2\hat{a}_z \end{aligned}$$

Magnitude of \vec{R}_{AB} :-

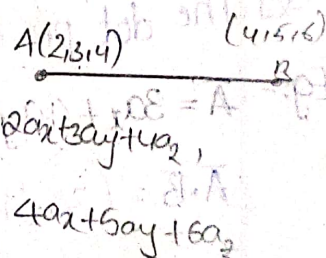
$$|\vec{R}_{AB}| = \sqrt{3^2 + 1^2 + 2^2} = \sqrt{9+1+4} = \sqrt{14}$$

unit vector:-

$$\hat{a}_{R_{AB}} = \frac{\vec{R}_{AB}}{|\vec{R}_{AB}|}$$

$$= \frac{3\hat{a}_x + \hat{a}_y + 2\hat{a}_z}{\sqrt{14}} = \frac{3\hat{a}_x + \hat{a}_y + 2\hat{a}_z}{\sqrt{14}}$$

unit vector $\hat{a}_{R_{AB}} = \frac{\vec{R}_{AB}}{|\vec{R}_{AB}|}$



$$= \frac{3a_x + a_y + 2a_z}{\sqrt{14}}$$

Vector addition and Substitution

$$\vec{A} + \vec{B} = \vec{B} + \vec{A}$$

$$\vec{A} + (\vec{B} + \vec{C}) = (\vec{A} + \vec{B}) + \vec{C}$$

$$\vec{A} - \vec{B} = \vec{B} - \vec{A}$$

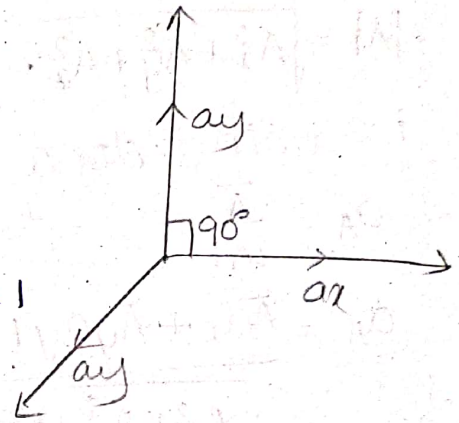
Scalar / dot product:-

$$A \cdot B = |A||B|\cos\theta$$

Eg:- $a_x \cdot a_x = |a_x||a_x|\cos(a_x \cdot a_x)$
 $= (1)(1)(1)$

similarly $a_y \cdot a_y = 1$ $a_z \cdot a_z = 1$

$a_x \cdot a_y = |a_x||a_y|\cos(a_x \cdot a_y)$
 $= (1)(1)(\cos 90^\circ)$



Similarly $a_x \cdot a_y = 0$; $a_y \cdot a_z = 0$; $a_z \cdot a_x = 0$

→ The dot product is scalar product there is no direction

Eg:- $\vec{A} = 3a_x + 4a_y + 5a_z$; $\vec{B} = 6a_x + 4a_y + 2a_z$

$$\vec{A} \cdot \vec{B} = (3a_x + 4a_y + 5a_z) \cdot (6a_x + 4a_y + 2a_z)$$

$$= 18 + 16 + 10$$

$$= 44$$

$$\therefore \vec{A} \cdot \vec{B} = \vec{B} \cdot \vec{A}$$

10-09-22 Cross product (or) Vector product:-

$$\vec{A} \times \vec{B} = |\vec{A}||\vec{B}|\sin\theta \hat{a}_n$$

Eg:- $a_x \times a_x = |a_x||a_x|\sin\theta \hat{a}_n$

$$= (1)(1)0$$

$$= 0$$

Similarly $a_y \times a_y = 0$; $a_z \times a_z = 0$.

$$a_1 \times a_2 = |a_1| |a_2| \sin \theta \hat{n}$$

$$= (1)(1)(1)a_2$$

$$= a_2$$

Similarly $a_2 \times a_1 = -a_1$; $a_1 \times a_1 = 0$

eg $\bar{A} \times \bar{B} = |\bar{B} \times \bar{A}|$

$$\bar{A} = 2a_1 + 3a_2 + 2a_3$$

$$\bar{B} = 3a_1 - 4a_2 - 2a_3$$

$$\bar{A} \times \bar{B} = (2a_1 + 3a_2 + 2a_3) \times (3a_1 - 4a_2 - 2a_3)$$

$$= -8a_2 - 4a_3 + 9a_3 - 6a_1 + 3a_2 - 4a_1$$

$$= -10a_1 - a_2 + a_3$$

$$|\bar{B} \times \bar{A}| = |(3a_1 - 4a_2 - 2a_3) \times (2a_1 + 3a_2 + 2a_3)|$$

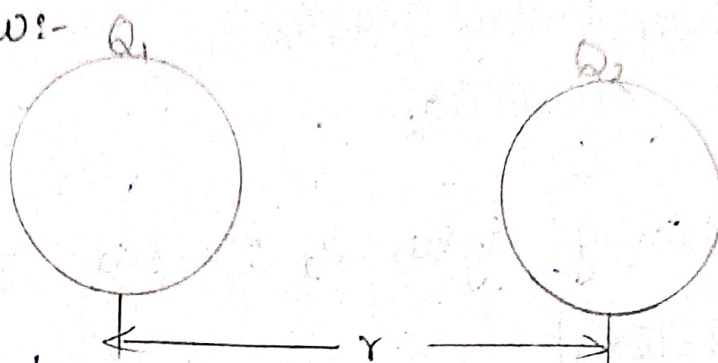
$$= |9a_3 + 3a_2 - 8a_2 - 4a_1 - 4a_3 - 6a_1|$$

$$= |-10a_1 - a_2 + a_3|$$

$$= |10a_1 + a_2 - a_3|$$

$$\therefore \bar{A} \times \bar{B} = |\bar{B} \times \bar{A}|$$

Columbs Law:-



- * consider the two point charges Q_1, Q_2 are separated by distance r . The charge Q_1 exerts force on Q_2 . While the charge Q_2 exerts a force on Q_1 .

The force exerted in between them is repulsive force if the charges are same polarity which is attractive if the charges are different polarities.

\therefore The force is directly proportional to the two charges Q_1 & Q_2

$$\therefore F \propto Q_1 \cdot Q_2$$

The force is inversely proportional to the square of distance between two charges

$$F \propto \frac{1}{r^2}$$

$$F \propto \frac{Q_1 Q_2}{r^2}$$

$$F = \frac{k Q_1 Q_2}{r^2}$$

$$F = \frac{Q_1 Q_2}{4\pi\epsilon r^2}$$

$$F = \frac{Q_1 Q_2}{4\pi\epsilon_0 r^2}$$

$$\therefore \mu = \mu_0 \mu_r$$

$$\mu_0 = 4\pi \times 10^{-7} \text{ H/m}$$

$$k = \frac{1}{4\pi\epsilon}$$

$$\epsilon = \epsilon_0 \epsilon_r$$

$$\epsilon_0 = 8.85 \times 10^{-12} \text{ F/m}$$

$$\epsilon_r = 1 \text{ (air / free space)}$$

Electric field intensity:

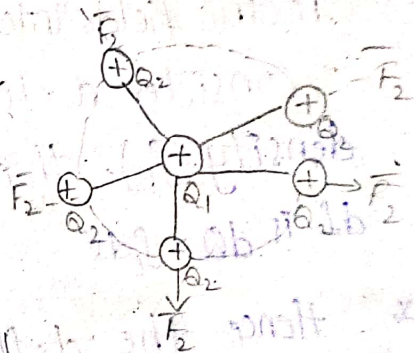
Consider a point charge Q_1 & another charge Q_2
 $\therefore Q_1$ is the point charge, Q_2 is the moved charge
 Thus there exists a region around a charge exerts a force on any other charge located in that region. is called electric field of that charge.

* The force exerted per unit charge is called electric field intensity or electric field strain.

* \therefore Force experienced by the charge Q_2 due to Q_1 is given by coulombs law.

$$F_Q = \frac{Q_1 Q_2}{4\pi\epsilon_0\epsilon_r R_{12}^2} \bar{a}_{12}$$

$$\bar{E} = \frac{Q_1}{4\pi\epsilon_0\epsilon_r R_{12}^2} \bar{a}_{12}$$



\therefore Electric field intensity

$$\bar{E} = \frac{F_2}{Q_2}$$

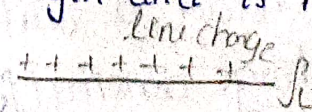
$$\therefore \bar{E} = \frac{Q_1}{4\pi\epsilon_0\epsilon_r R_{12}^2} \bar{a}_{12} \text{ N/C} = \frac{\text{Newton's}}{\text{Coulombs}}$$

Line charge:-

The charge may be spreaded all along a line which may be finite or infinite. Such charge uniformly distributed along a line is called line charge.

\therefore The charge density of the line charge is denoted by ρ_L and defined as charge per unit length and is measured in coulombs per meter.

$$\therefore \rho_L = \frac{\text{total charge in coulombs}}{\text{Total length in meters}}$$

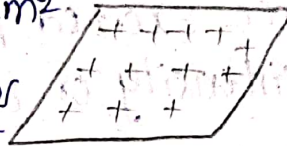


$$Q = \int dq \quad dq = \rho_L dl$$

Surface charge

If the charge is distributed uniformly over two dimensional surface then it is called a surface charge or sheet of charge. It is denoted as ρ_s & is defined as the charge per unit surface area & it measure in coulombs per meter² C/m^2 .

$$\rho_s = \frac{\text{Total charge in coulombs}}{\text{Total area in } m^2}$$



Electric field intensity due to line charge

Consider a line charge distribution having a charge density (ρ_L), that charge dq on the differential length dl is $dQ = \rho_L dl$

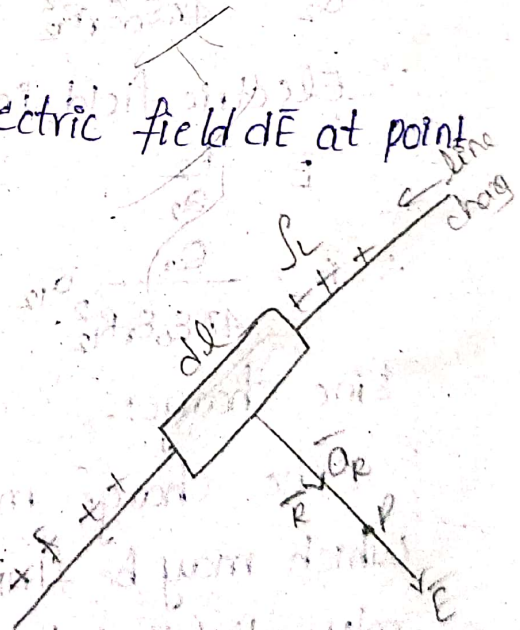
* Hence the differential electric field $d\vec{E}$ at point P due to dq is given by

$$d\vec{E} = \frac{dq}{4\pi\epsilon_0 R^2} \vec{a}_R$$

$$d\vec{E} = \frac{\rho_L dl}{4\pi\epsilon_0 R^2} \vec{a}_R$$

$$\int d\vec{E} = \int \frac{\rho_L dl}{4\pi\epsilon_0 R^2} \vec{a}_R$$

$$\vec{E} = \int \frac{\rho_L dl}{4\pi\epsilon_0 R^2} \vec{a}_R$$



Electric field intensity due to surface charge

Consider a surface charge distribution, having a charge density (ρ_s). the charge dq on the differential area ds is

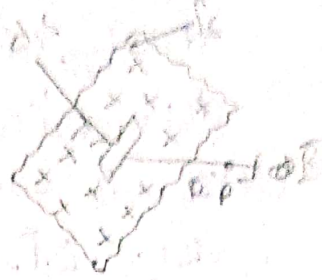
$$dq = \rho_s ds$$

* Hence differential equation Electrical-field $d\vec{E}$ at a point P due to dQ is given by

$$d\vec{E} = \frac{dQ}{4\pi\epsilon_0 R^2} \vec{a}_R$$

$$\therefore dQ = \int_S ds$$

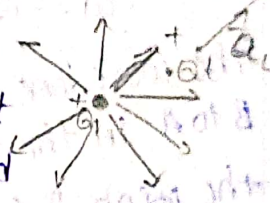
$$d\vec{E} = \int \frac{\rho ds}{4\pi\epsilon_0 R^2} \vec{a}_R$$



Hence the total \vec{E} at a point P is due to obtained by integrating \vec{E} over the surface area S on which charge is distribution $\therefore \vec{E} = \int \frac{\rho ds}{4\pi\epsilon_0 R^2} \vec{a}_R$.

22-01-22 Workdone in moving a point charge in electric field

* Consider a positive charge Q_1 & its electric field \vec{E} , if a positive test charge Q_2 is placed in the field it will move due to force of repulsion



* Let the moment of charge Q_2 is dL and denoted by unit vector \vec{a}_L

* i.e according Coulombs law the force exerted by the field \vec{E} is given by $\vec{F} = \frac{F}{Q_2}$

$$\vec{F} = Q_2 \vec{E} \text{ Newtons}$$

* The component of \vec{F} in the direction unit vector \vec{a}_L is given by $\vec{F}_L = \vec{F} \cdot \vec{a}_L$

$$\vec{F}_L = Q_2 \vec{E} \cdot \vec{a}_L$$

* i.e It is necessary to apply the force which is equal to & opposite to the force exerted by the field in the direction in the dL .

$$\vec{F}_{\text{applied}} = -\vec{F}_L$$

$$\vec{F}_{\text{app}} = -Q_2 \vec{E} \cdot \vec{a}_L$$

* \therefore differential workdone by an external source in moving the charge Q_2 through a distance dL

against the dielectric field ϵ is given by $w = k \epsilon$

differentiate

$$dw = f_{app} \times dd$$

$$dw = -Q_t \bar{E} \bar{a}_L \times dd$$

Since $dL \times \bar{a}_L = d\vec{L}$ \longrightarrow using dot vector

$$dw = -Q_t \bar{E} \cdot d\vec{L}$$

$$\therefore \int dw = \int -Q_t \bar{E} \cdot d\vec{L}$$

$$w = - \int_{\text{initial}}^{\text{final}} Q_t \bar{E} \cdot d\vec{L}$$

Potential difference

The workdone in moving a point charge Q from point B to A in the electric field \bar{E} is given by

$$w = -Q \int_B^A \bar{E} \cdot d\vec{L}$$

The workdone per unit charge in moving unit charge from B to A in the field \bar{E} is called potential difference between the points B & A and it is denoted by V and it is measured in joules/coulombs

$$\text{potential difference} = \frac{\text{workdone}}{\text{charge}}$$

$$= - \frac{Q \int_B^A \bar{E} \cdot d\vec{L}}{Q}$$

$$= - \int_B^A \bar{E} \cdot d\vec{L}$$

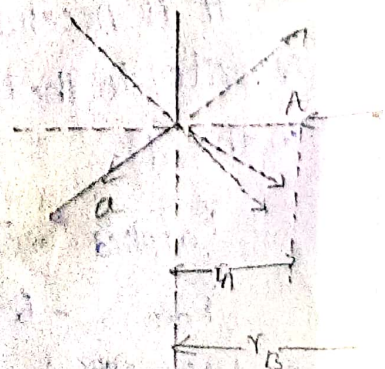
Potential due to point charge:

* consider a point charge located at the origin of a spherical coordinate system, producing \bar{E} radially in all directions

* Assuming the free space, the field \bar{E} due to a point charge Q at a point having radial distance r from origin is given by

$$\bar{E} = \frac{Q}{4\pi\epsilon_0 r^2} \bar{a}_r$$

The differential line is in spherical system is



$$V_{AB} = -\int E \cdot dl$$

$$V_{AB} = -\int_B^A \frac{Q}{4\pi\epsilon_0 r^2} \bar{a}_r (dr \bar{a}_r + r d\theta \bar{a}_\theta + r \sin\theta d\phi \bar{a}_\phi)$$

$$V_{AB} = -\int_B^A \frac{Q}{4\pi\epsilon_0 r^2} dr = \frac{Q}{4\pi\epsilon_0} \int_B^A r^{-2} dr = \frac{-Q}{4\pi\epsilon_0} \left[\frac{1}{r} \right]_B^A = \frac{+Q}{4\pi\epsilon_0} \left[\frac{1}{A} - \frac{1}{B} \right]$$

$$V_{AB} = \frac{Q}{4\pi\epsilon_0} \left[\frac{1}{r_A} - \frac{1}{r_B} \right] \text{ volts}$$

Potential gradient:-

The rate of change of potential with respect to distance is called potential gradient

$$\therefore \text{potential } v = -\int E \cdot dl = \frac{Q}{4\pi\epsilon_0 r}$$

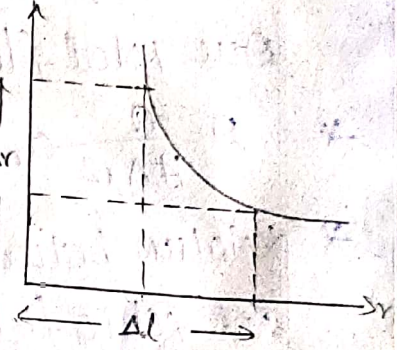
The potential decreases as distances of or point from the charge increases. From an elementary length Δl is given by V_{AB}

$$V_{AB} = \Delta V = -E \cdot \Delta l$$

\therefore Rate of change of potential difference

$$\frac{dv}{dl} = \lim_{\Delta l \rightarrow 0} \frac{\Delta V}{\Delta l} = \text{potential gradient}$$

The potential gradient is nothing but the slope of the graph of potential against distance at a point where elementary length is considered.



2 marks

The total number of electric lines of force or flux in any particular electric field is called Electric flux

Properties of flux lines:-

- *
- *
- *
- *

The flux lines starts from positive charge and terminate on the negative charge

There are more number of lines, i.e. (thru) crowding of lines if electric field is stronger

These lines are parallel and never cross each other

The lines always enter or leave the charged surface normally

electric flux density:-

The total electric flux passing through the unit

unit surface area is called the electric flux density

* It is denoted by \bar{D} and is measured in coulombs/m².

$$\bar{D} = \frac{\psi}{S}$$

* where ψ = Total electric flux in coulombs

$$S = \text{Total Surface Area in meter}^2 (\text{m}^2) = 4\pi r^2 \text{ of sphere.}$$

$$\therefore \bar{D} = \frac{Q}{4\pi r^2}$$

Due Total flux = charge

$$\bar{D} = \frac{Q}{4\pi r^2} \bar{a}_r \text{ in vector form}$$

* Relation between \bar{D} & \bar{E}

$$\therefore \text{Electric field intensity } \bar{E} = \frac{Q}{4\pi\epsilon_0 r^2} \bar{a}_r \rightarrow (1)$$

$$\text{Electric flux density } \bar{D} = \frac{Q}{4\pi r^2} \bar{a}_r \rightarrow (2)$$

dividing (2) with (1)

$$\frac{(2)}{(1)} \Rightarrow \frac{\bar{D}}{\bar{E}} = \frac{\frac{Q}{4\pi r^2} \bar{a}_r}{\frac{Q}{4\pi\epsilon_0 r^2} \bar{a}_r}$$

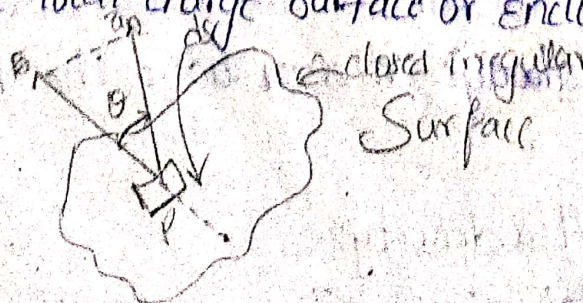
$$\frac{\bar{D}}{\bar{E}} = \frac{1}{\epsilon_0}$$

$$\frac{\bar{D}}{\bar{E}} = \epsilon_0$$

$$\bar{D} = \bar{E} \epsilon_0$$

Gauss's law:-

The electric flux passing through any closed surface is equal to the total charge surface or enclosed by the surface



* The total surface enclosed by the irregular closed surface is Q in coulombs.

* Consider a small differential surface ds at a point P , the direction of \vec{E} as well as its magnitude is going to change from the point to point on the surface.

* $\therefore ds = an ds$ ($\because an =$ normal to the surface ds at point P)

* We already know that

$$D = \frac{\psi}{s}$$

$$d\psi = \vec{D} \cdot d\vec{s}$$

* Since $D_n = D \cos \theta =$ component of the \vec{D} in the direction of normal surface ds .

* $\therefore d\psi = D_n \cdot ds$

$$D_n = D \cos \theta$$

$$d\psi = D \cos \theta \cdot ds$$

$$d\psi = D \cdot ds$$

$$\int d\psi = \oint \vec{D} \cdot d\vec{s}$$

$$\psi = Q$$

$$\therefore \psi = \int d\psi = \oint \vec{D} \cdot d\vec{s} = Q$$

28-11-22
* If there is a line charge with line charge density (ρ_L), then $\psi = Q = \int \rho_L dl$

* If there is a surface charge with surface charge density (ρ_s) then $\psi = Q = \int \rho_s ds$

* If there is a volume charge with volume charge density ρ_v then $\psi = Q = \int \rho_v dV$

Divergence in different coordinate system

- * There are classified into three types
 - i. Cartesian / rectangular system
 - ii. cylindrical co-ordinate system
 - iii. spherical co-ordinate system

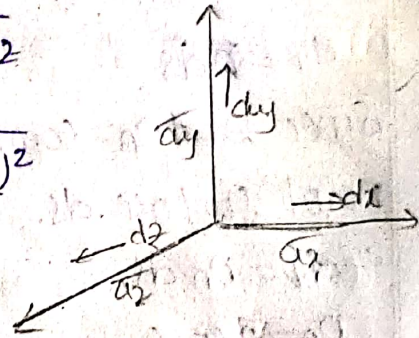
i. Cartesian / Rectangular system:

* consider a point (x, y, z) in rectangular co-ordinate system and $d\vec{L}$ is the differential length

$$\therefore d\vec{L} = dx\vec{a}_x + dy\vec{a}_y + dz\vec{a}_z$$

$$|d\vec{L}| = \sqrt{(dx)^2 + (dy)^2 + (dz)^2}$$

$$\Delta \cdot \vec{A} = \text{div} \vec{A} = \frac{\partial A_x}{\partial x} + \frac{\partial A_y}{\partial y} + \frac{\partial A_z}{\partial z}$$



ii. cylindrical co-ordinate system:-

* consider a point (r, ϕ, z) in cylindrical co-ordinate system. dr is the differential length in r direction, $r d\phi$ is the differential length in ϕ direction, dz is the differential length in z direction

$$d\vec{L} = dr\vec{a}_r + r d\phi\vec{a}_\phi + dz\vec{a}_z$$

$$|d\vec{L}| = \sqrt{dr^2 + (r d\phi)^2 + (dz)^2}$$

$$\Delta \cdot \vec{A} = \text{div} \vec{A} = \frac{1}{r} \frac{\partial}{\partial r} (r A_r) + \frac{1}{r} \frac{\partial A_\phi}{\partial \phi} + \frac{\partial A_z}{\partial z}$$

iii. Spherical co-ordinate system:-

Consider a point (r, θ, ϕ) in spherical co-ordinate system. $d\vec{L}$ is the differential length

$$\therefore d\vec{L} = dr\vec{a}_r + r d\theta\vec{a}_\theta + r \sin\theta d\phi\vec{a}_\phi$$

$$|d\vec{L}| = \sqrt{(dr)^2 + (r d\theta)^2 + (r \sin\theta d\phi)^2}$$

$$\Delta \cdot \vec{A} = \text{div} \vec{A} = \frac{1}{r^2} (r^2 A_r) + \frac{1}{r \sin\theta} \frac{\partial}{\partial \theta} (\sin\theta A_\theta) + \frac{1}{r \sin\theta} \frac{\partial A_\phi}{\partial \phi}$$

Maxwell's first equation:

According to Gauss's law

$$\Phi = Q = \oint \vec{D} \cdot d\vec{S} \rightarrow (1)$$

Both side dividing per unit volume ΔV

$$\frac{Q}{\Delta V} = \frac{\oint \vec{D} \cdot d\vec{S}}{\Delta V}$$

Taking limits $\Delta V \rightarrow 0$

$$\text{i.e. } \lim_{\Delta V \rightarrow 0} \frac{Q}{\Delta V} = \lim_{\Delta V \rightarrow 0} \frac{\oint \vec{D} \cdot d\vec{S}}{\Delta V} \rightarrow (2)$$

\therefore The divergence of electric flux density \vec{D} is given by

$$\Delta \cdot \vec{D} = \text{div} \vec{D} = \frac{\partial D_x}{\partial x} + \frac{\partial D_y}{\partial y} + \frac{\partial D_z}{\partial z}$$

$$\therefore \text{div} \vec{D} = \lim_{\Delta V \rightarrow 0} \frac{\oint \vec{D} \cdot d\vec{S}}{\Delta V} \quad (\because \lim_{\Delta V \rightarrow 0} \frac{Q}{\Delta V} = \rho)$$

from (2) $\rho_v = \text{div} \vec{D}$

$$\therefore \text{div} \vec{D} = \rho_v$$

$$\Delta \cdot \vec{D} = \rho_v$$

This is the volume charge density around a point the above equation is called Maxwell's first equation / point form of Gauss law or Gauss law in differential form

Poisson's & Laplace equations:

We know that Gauss law in the point form

$$\Delta \cdot \vec{D} = \rho_v \rightarrow (1)$$

where $\rho_v =$ volume charge density, $\vec{D} =$

$\vec{D} =$ Electric flux density

$$\therefore \vec{D} = \epsilon \vec{E} \rightarrow (2)$$

eq (2) in eq (1)



$$\Delta \cdot \epsilon \vec{E} = \rho_v \rightarrow (3)$$

\therefore electric potential gradient $\vec{E} = -\Delta V$

$$\Delta \cdot \epsilon (-\Delta V) = \rho_v$$

$$\epsilon \Delta \cdot \Delta V = -\rho_v$$

$$\Delta^2 V = \frac{-\rho_v}{\epsilon} \rightarrow (4)$$

The above equation is called Poisson's equation. In certain region volume charge density is zero ($\rho_v = 0$), which is true for dielectric medium in the Poisson's equation then the

$$\Delta^2 V = 0$$

The above equation is called Laplace equation, the Δ^2 operation is called the Laplace of V .

Δ^2 operation in different co-ordinate system:

* The potential (V) can be expressed in 3-coordinates systems $V(x, y, z)$, $V(r, \phi, z)$, $V(r, \theta, \phi)$ depending upon Δ^2 operation required for Laplace equation must be used.

* Cartesian coordinate system:-

$$\Delta V = \frac{\partial V}{\partial x} \bar{a}_x + \frac{\partial V}{\partial y} \bar{a}_y + \frac{\partial V}{\partial z} \bar{a}_z$$

$$\Delta \cdot \Delta V = \frac{\partial}{\partial x} \left(\frac{\partial V}{\partial x} \right) + \frac{\partial}{\partial y} \left(\frac{\partial V}{\partial y} \right) + \frac{\partial}{\partial z} \left(\frac{\partial V}{\partial z} \right)$$

$$\Delta^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

* Cylindrical co-ordinate system:-

$$\Delta^2 V = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial V}{\partial r} \right) + \frac{1}{r^2} \left(\frac{\partial^2 V}{\partial \phi^2} \right) + \left(\frac{\partial^2 V}{\partial z^2} \right) = 0$$

* Spherical co-ordinate system:-

$$\Delta^2 V = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V}{\partial \phi^2}$$

pb:1 Determine wheather or not the potential fields satisfy the Laplace equations

(i) $V = x^2 - y^2 + z^2$ (ii) $V = r \cos \phi + z$ (iii) $r \cos \theta + \phi$

Sol given $V = x^2 - y^2 + z^2$

$\Delta^2 V = ?$

$\Delta V = \frac{\partial V}{\partial x} + \frac{\partial V}{\partial y} + \frac{\partial V}{\partial z}$

$= 2x - 2y + 2z$

$\Delta V = 2(x - y + z)$

$\Delta^2 V = 2 \left(\frac{\partial}{\partial x} (x - y + z) + \frac{\partial}{\partial y} (x - y + z) + \frac{\partial}{\partial z} (x - y + z) \right)$

$= 2(1 - 1 + 1)$

$\Delta^2 V = 2 \neq 0$

(ii) $V = r \cos \phi + z$

$\Delta^2 V = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} (r \cos \phi + z) \right) + \frac{1}{r^2} \left(\frac{\partial^2 (r \cos \phi + z)}{\partial \phi^2} \right) + \frac{\partial^2}{\partial z^2} (r \cos \phi + z)$

$= \frac{1}{r} \frac{\partial}{\partial r} (r \cos \phi) + \frac{1}{r^2} \left(\frac{\partial}{\partial \phi} (r) (-\sin \phi) \right) + \frac{\partial^2}{\partial z^2} (1)$

$= \frac{1}{r} \cos \phi + \left(-\frac{1}{r^2} \right) r \cos \phi + 0$

$= \frac{1}{r} (\cos \phi - \frac{1}{r} \cos \phi)$

$= \frac{\cos \phi}{r} (1 - \frac{1}{r}) = 0$

(iii) given $V = r \cos \theta + \phi$

$\Delta^2 V = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V}{\partial \phi^2}$

$= \frac{1}{r^2} \left(\frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} (r \cos \theta + \phi) \right) \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} (r \cos \theta + \phi) \right)$

$+ \frac{1}{r^2 \sin^2 \theta} \left(\frac{\partial}{\partial \phi} \left(\frac{\partial}{\partial \phi} (r \cos \theta + \phi) \right) \right)$

$$\Delta^2 u = \frac{1}{r^2} \left[\frac{\partial}{\partial r} [r^2 \cos \theta] \right] + \frac{1}{r^2 \sin \theta} \left[\frac{\partial}{\partial \theta} \sin \theta \left(\frac{\partial}{\partial \phi} \sin \theta \right) \right]$$

$$+ \frac{1}{r^2 \sin^2 \theta} \left[\frac{\partial}{\partial \phi} (1) \right]$$

$$\Delta^2 u = \frac{1}{r^2} 2r \cos \theta + \frac{-1 \sin \theta}{r^2 \sin \theta} \left[\frac{\partial}{\partial \theta} r \sin^2 \theta \right]$$

$$+ \frac{1}{r^2 \sin^2 \theta} (0)$$

$$\Delta^2 u = \frac{2}{r} \cos \theta - \frac{r}{r^2 \sin^2 \theta} (2 \sin \theta \cos \theta) + 0$$

$$= \frac{2}{r} \cos \theta - \frac{1}{r \sin^2 \theta} 2 \sin \theta \cos \theta$$

$$= \frac{2}{r} \cos \theta - \frac{2 \cos \theta}{r \sin \theta}$$

$$= \frac{2}{r} \cos \theta \left(1 - \frac{1}{\sin \theta} \right)$$

$$= 0$$

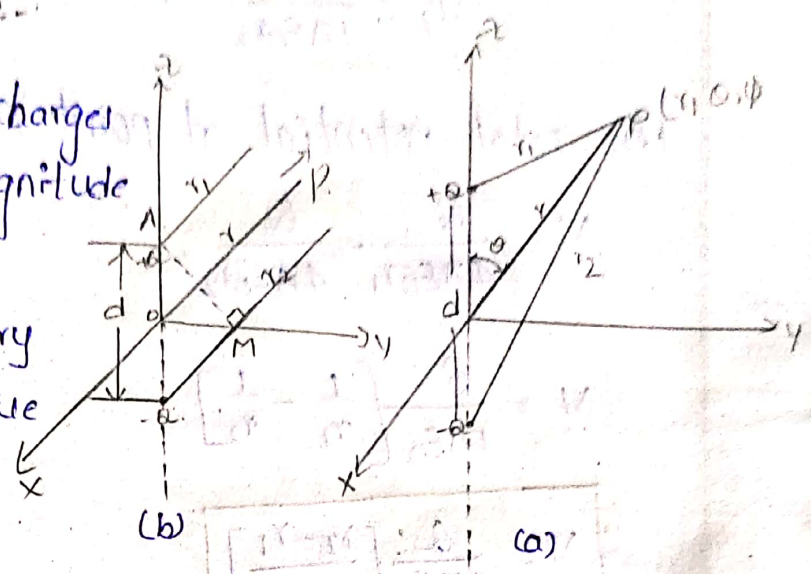
12/10/22

UNIT - II

ELECTRIC DIPOLE

Electric dipoles:-

The two point charges of a equal magnitude but opposite sign Separated by very small distance give rise to an electric dipole.



consider an electric dipole the two point charges +Q, -Q. are Separated by very small distance d.

consider a point P(r, theta, phi) in spherical co-ordinate system

Let "O" be the mid point of "AB". The distance of point "P" from "A" is "r1", while the distance of point "P" from "B" is "r2". The distance of point "P" from point "O" is "r".

To find "E", we will find out potential "v" at point "P" due to an electric dipole. Then using

$$\vec{E} = -\Delta v \quad | \quad \vec{E} = -\nabla v$$

In Spherical co-ordinates the potential at point "P" due to the charge "+Q" is given by

$$v_1 = \frac{+Q}{4\pi\epsilon_0 r_1}$$

At potential at point "p" due to the charge " $-Q$ " is given by

$$V_2 = \frac{-Q}{4\pi\epsilon_0 r_2}$$

The total potential at point "p" $V = V_1 + V_2$

$$V = \frac{Q}{4\pi\epsilon_0 r_1} - \frac{Q}{4\pi\epsilon_0 r_2}$$

$$V = \frac{Q}{4\pi\epsilon_0} \left[\frac{1}{r_1} - \frac{1}{r_2} \right]$$

$$V = \frac{Q}{4\pi\epsilon_0} \left[\frac{r_2 - r_1}{r_1 r_2} \right]$$

The r_1, r_2, r are assume to be parallel to each other. AM is drawn perpendicular from "A" on " r_2 "

$$PB = BM + PM \quad \text{from fig(b)}$$

$$BM = PB - PM$$

$$PM = PA$$

$$BM = r_2 - r_1$$

$$(r_2 - r_1) = d \cos \theta \quad \text{[} \cdot \text{ BM} = d \cos \theta \text{]}$$

$$\text{Since } r_1 = r_2 = r$$

$$V = \frac{Q}{4\pi\epsilon_0} \left[\frac{d \cos \theta}{r^2} \right]$$

$$\vec{E} = -\nabla V$$

$$\vec{E} = -\nabla V = - \left[\frac{\partial V}{\partial r} \vec{a}_r + \frac{1}{r} \frac{\partial V}{\partial \theta} \vec{a}_\theta + \frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi} \vec{a}_\phi \right]$$

$$= - \left[\frac{Q d \cos \theta}{4\pi\epsilon_0} \frac{\partial}{\partial r} \left(\frac{1}{r^2} \right) \vec{a}_r + \frac{Q d}{4\pi\epsilon_0} \frac{\partial}{\partial \theta} \cos \theta \vec{a}_\theta + \frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi} \vec{a}_\phi \right]$$

$$= - \left[\frac{Q d \cos \theta}{4\pi\epsilon_0} \left(\frac{-2}{r^3} \right) \vec{a}_r + \frac{Q d}{4\pi\epsilon_0} (-\sin \theta) \vec{a}_\theta \right]$$

$$= + \frac{Qd \cos \theta}{4\pi\epsilon_0 r^3} \bar{a}_r + \frac{Qd \sin \theta}{4\pi\epsilon_0 r^3} \bar{a}_\theta$$

$$\bar{E} = -\nabla V = \frac{Qd}{4\pi\epsilon_0 r^3} [2\cos\theta \bar{a}_r + \sin\theta \bar{a}_\theta]$$

13/10/22 Dipole moment:-

The product of charge and distance is called dipole moment and is denoted by "P" and is measured in coulombs meters (C-m).

$$\therefore \bar{P} = Q\bar{d}$$

The dipole moment is measured in coulombs meter

$$\text{Now } \bar{P} \cdot \bar{a}_r = Q\bar{d} \cdot \bar{a}_r$$

$$\bar{P} \cdot \bar{a}_r = Q(\bar{d} \cdot \bar{a}_r)$$

$$\bar{P} \cdot \bar{a}_r = Qd \|\bar{a}_r\| \cos \theta$$

$$\bar{P} \cdot \bar{a}_r = Qd \cos \theta$$

$$\bar{P} \cdot \bar{a}_r = Qd \cos \theta$$

Hence the expression of potential "V" can be expressed

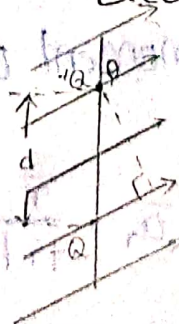
$$\text{as } V = \frac{Qd \cos \theta}{4\pi\epsilon_0 r^2}$$

$$V = \frac{\bar{P} \cdot \bar{a}_r}{4\pi\epsilon_0 r^2}$$

$$\bar{E} = \frac{Qd}{4\pi\epsilon_0 r^3} [2\cos\theta \bar{a}_r + \sin\theta \bar{a}_\theta]$$

Torque on electric dipole in an electric field:-

Consider an electric dipole in an uniform electric field \bar{E} . Making an angle θ w.r.t dipole axis



The charges "+Q" & "-Q" experienced a force due to an electric field \vec{E} is equal to magnitude but opposite direction.

* Torque = Force \times displacement
Or

Torque = force \times perpendicular distance of a separation.

$$\tau = F \times l$$

$$\tau = Eq \times l$$

$$\tau = EqL$$

$$E = \frac{F}{Q}$$

$$F = Eq$$

Where L = perpendicular distance of separation between the two forces

$$\sin \theta = \frac{l}{d}$$

$$l = d \sin \theta$$

$$\tau = Eqd \sin \theta$$

While $\sin \theta$ is the cross product b/w the two the torque on the dipole is expressed as

$$\tau = \vec{p} \times \vec{E}$$

↳ cross

pb: 1 A dipole having moment $\vec{p} = 3a_1 - 5a_2 + 10a_3$ Ncm is locate at $Q(1, 2, -4)$ in free space. Find V at $P(2, 3, 4)$

dipole moment of potential $V = \frac{\vec{p} \cdot \vec{a}_r}{4\pi\epsilon_0 r^2}$

$$\vec{a}_r = ? \quad \vec{a}_r = \frac{\vec{r}}{|\vec{r}|}$$

$$\vec{r} = (2-1)\vec{a}_1 + (3-2)\vec{a}_2 + (4+4)\vec{a}_3$$

$$r = a_x + a_y + 8a_z$$

$$|r| = \sqrt{1^2 + 1^2 + 8^2}$$

$$|r| = \sqrt{66} = 8.1$$

$$\bar{a}_r = \frac{r}{|r|} = \frac{a_x + a_y + 8a_z}{\sqrt{66}}$$

$$\bar{P} \cdot \bar{a}_r = (3a_x - 5a_y + 10a_z) \cdot \frac{a_x + a_y + 8a_z}{\sqrt{66}}$$

$$= \frac{3 - 5 + 80}{\sqrt{66}}$$

$$= \frac{78}{\sqrt{66}} \times 10^{-9} \text{ cm}$$

$$= 9.60 \times 10^{-9} \text{ cm}$$

$$V = \frac{9.60 \times 10^{-9}}{4\pi \times 8.854 \times 10^{-12} (\sqrt{66})^2}$$

$$V = 1.30 \text{ volts}$$

prob. 2 Compute the Torque for a dipole consisting of 1 micro columbia charges in the electric field $\bar{E} = 10^3(2a_x - a_y - a_z) \text{ V/m}$ Separated by 1mm & located on the z axis (only) at the origin.

Given

$$d = 1 \text{ mm}$$

$$d = 1 \times 10^{-3} \text{ m}$$

$$Q = 1 \text{ mc}$$

$$Q = 10^{-6} \text{ coulomb}$$

$$\bar{E} = 10^3(2a_x - a_y - a_z)$$

$$\tau = \bar{P} \times \bar{E}$$

$$\bar{P} = Qd$$

$$\bar{P} = 10^{-6} \times 10^{-3}$$

$$\bar{P} = 10^{-9}$$

$$\bar{P} = 1 \times 10^{-9} a_z \text{ (due to z-axis)}$$

$$\gamma = 10^{-9} a_2 \times (10^3 (2\bar{a}_x - \bar{a}_y - \bar{a}_z))$$

$$= 10^{-9} a_2 \times (10^3 2\bar{a}_x - 10^3 \bar{a}_y - 10^3 \bar{a}_z)$$

$$x = 10^{-6} (10^{-9} \times 10^3) + (10^{-9} \times 10^3)$$

$$x = 10^{-6} 2\bar{a}_x - 10^{-6} \bar{a}_z$$

$$v = 10^{-6} (2 - 1)$$

$$= 2 \cdot 10^{-6} \bar{a}_x - 10^{-6} \bar{a}_z$$

15/10/22 polarization:-

Consider an atom of a dielectric is consist of a nucleus with positive charge and negative charge in the form of revolving electrons in the orbits

The negative charge is thus consider to be in the form of cloud of electrons.

The number of positive charge is the same as negative charge and hence atom is electrically neutral

Both positive and negative charges can be assume to be point charges of equal amount coinciding at the centre

hence there cannot exist an electric dipole this is called unpolarized atom

The separation between the nucleus and the centre of the electron cloud, such an atom is called polarized atom.

The dipole gets aligned with the applied field. This process is called polarization of dielectrics. There are two types of dielectric

1. Non-polar:-

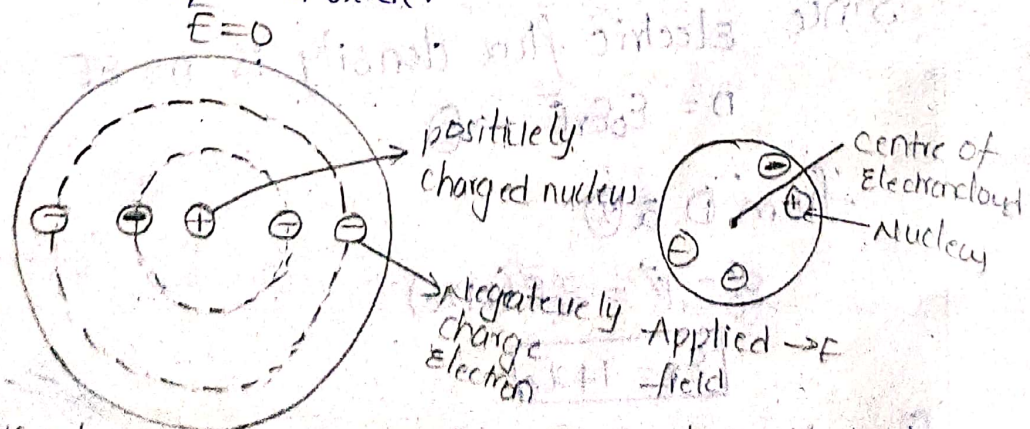
In non-polar molecules, the dipole arrangement is totally absent. In the absence of electric field E , it results only when an externally field E is applied to it

2. Polar:-

In polar molecules, the permanent displacements between the centres of positive and negative charges exists thus dipole arrangements exists without application of E . The dipole experienced Torque and then align with the direction of the applied field E . This is called polarization of polar molecules.

for example of non-polar molecules are Hydrogen, Oxygen and rare gases.

The Example of polar molecules are water, hydrochloric acid Sulphur dioxide.



unpolarized atom of a dielectric polarized atom

Mathematical expression for polarization:

dipole moment $p = qd$

The total dipole moment is to be obtained using

Superposition principle is

$$P_{\text{total}} = Q_1 \bar{d}_1 + Q_2 \bar{d}_2 + Q_3 \bar{d}_3 + \dots + Q_n \bar{d}_n = \sum_{i=1}^{n\Delta V} Q_i \bar{d}_i$$

$$\bar{P} = \lim_{\Delta V \rightarrow 0} \frac{\sum_{i=1}^{n\Delta V} Q_i \bar{d}_i}{\Delta V}$$

The polarization \bar{P} is defined as the total dipole moment per unit volume. It is measured in Coulombs per square meters C/m^2 .

$$\bar{P} = \lim_{\Delta V \rightarrow 0} \frac{\sum_{i=1}^{n\Delta V} Q_i \bar{d}_i}{\Delta V}$$

flux density in dielectric is $\bar{D} = \epsilon_0 \bar{E} + \bar{P}$
for isotropic and non-linear medium \bar{P} & \bar{E} are parallel to each other

$$\therefore \bar{P} = \chi_e \epsilon_0 \bar{E}$$

where χ_e = dimensionless quantity called electric susceptibility of the material

$$\bar{D} = \epsilon_0 \bar{E} + \chi_e \epsilon_0 \bar{E}$$

$$\bar{D} = \epsilon_0 \bar{E} (\chi_e + 1) \rightarrow \text{①}$$

Since electric flux density is $D = \epsilon E$

$$D = \epsilon_0 \epsilon_r \bar{E} \rightarrow \text{②}$$

from ① & ②

$$\epsilon_0 \epsilon_r \bar{E} = \epsilon_0 \bar{E} (1 + \chi_e)$$

$$\boxed{\epsilon_r = 1 + \chi_e}$$

$(1 + \chi_e)$ is defined as relative permittivity or the dielectric constant of the dielectric material

21/10/22
Boundary conditions:

- * When an electric field passes from one medium to another medium
- * The conditions existing at the boundary of the two media when field passes from one medium to other are called boundary conditions

It is depending upon the nature of the media there are two boundary conditions

1. Boundary between conductor and free space (or) boundary between conductor and dielectric
2. Boundary between two dielectrics with different properties

- * For studying the boundary conditions the max wells eqns for electrostatics are required

$$\oint \vec{E} \cdot d\vec{l} = 0 \quad \& \quad \oint \vec{D} \cdot d\vec{s} = Q$$

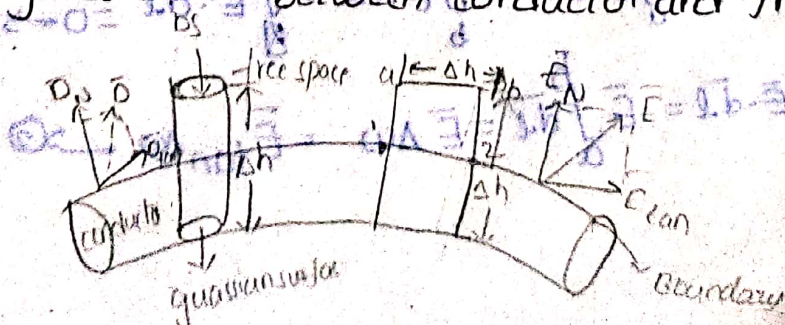
- * The field intensity \vec{E} is required to be decomposed into two components. They are

1. Tangentially to the boundary (\vec{E}_{tan}) E_t
2. Normal to the boundary (\vec{E}_N)

$$\vec{E} = \vec{E}_{tan} + \vec{E}_N$$

Similarly decomposition is require for flux density (\vec{D})

Boundary condition between conductor and free space:



Consider a boundary between conductor and free space. The conductor is ideal having infinite conductivity. Such conductors are copper, silver...

*] For ideal conductor it is known that

1. The field intensity inside a conductor is zero and flux density inside a conductor is zero
2. No charge can exist within a conductor. The charge appears on the surface in the form of surface charge density
3. The charge density within the conductor is zero
4. To determine the boundary conditions let us use a closed path and Gaussian surface
5. Let \vec{E} electric field intensity can be resolved into two components
 - i. The component tangentially to the surface (\vec{E}_{tan})
 - ii. The component normal to the surface (\vec{E}_N)

Consider a rectangular closed path $abcd$. We know that enclosed surface

$$\textcircled{1} \quad \oint \vec{E} \cdot d\vec{l} = 0$$

$$\oint \vec{E} \cdot d\vec{l} = \int_a^b \vec{E} \cdot d\vec{l} + \int_b^c \vec{E} \cdot d\vec{l} + \int_c^d \vec{E} \cdot d\vec{l} + \int_d^a \vec{E} \cdot d\vec{l} = 0 \quad \textcircled{1}$$

$$\oint \vec{E} \cdot d\vec{l} = \int_a^b \vec{E} \cdot d\vec{l} + \int_b^c \vec{E} \cdot d\vec{l} + \int_c^d \vec{E} \cdot d\vec{l} = 0 \rightarrow \textcircled{1}$$

$$\int_a^b \vec{E} \cdot d\vec{l} = E \int_a^b d\vec{l} = E \Delta w = \vec{E}_{tan} \Delta w \rightarrow \textcircled{2}$$

$$\int_b^c \vec{E} \cdot d\vec{l} = \int_b^a \vec{E} \cdot d\vec{l} + \int_a^c \vec{E} \cdot d\vec{l} \rightarrow \text{path closed} \therefore$$

$$\int_b^c \vec{E} \cdot d\vec{l} = \vec{E} \frac{\Delta h}{2} = \vec{E}_N \frac{\Delta h}{2} \rightarrow \textcircled{3}$$

$$\int_a^b \vec{E} \cdot d\vec{l} = -\vec{E}_N \frac{\Delta h}{2} \rightarrow \textcircled{4}$$

$$\vec{E}_{\tan} \Delta w + \vec{E}_N \frac{\Delta h}{2} - \vec{E}_N \frac{\Delta h}{2} = 0$$

$$\vec{E}_{\tan} \Delta w = 0$$

$$\therefore \vec{E}_{\tan} = 0$$

$$\therefore \vec{D}_{\tan} = 0$$

Tangentially component of electric field intensity and electric flux density is zero at the boundary b/w conductor and free space

According to Gauss law $\oint \vec{D} \cdot d\vec{s} = Q$ To find the form of right circular cylinder

The surface integral must be calculated top & bottom and lateral

$$\int_{\text{top}} \vec{D} \cdot d\vec{s} + \int_{\text{bottom}} \vec{D} \cdot d\vec{s} + \int_{\text{lateral}} \vec{D} \cdot d\vec{s} = Q$$

$$\int_{\text{top}} \vec{D} \cdot d\vec{s} = Q$$

$$\vec{D} \int_{\text{top}} d\vec{s} = Q$$

$$\vec{D}_N \Delta s = Q \rightarrow \textcircled{1}$$

\therefore Surface charge intensity $\rho_s \Delta s = \frac{Q}{\Delta s}$ (1)

comparing Eq (1) & (2)

$$\bar{D}_N \Delta s = \rho_s \Delta s$$

$$\bar{D}_N = \rho_s$$

Normal component of the Electric flux density

$$\bar{D}_N = \epsilon \bar{E}_N$$

$$\bar{D}_N = \epsilon_0 \bar{E}_N$$

$$\rho_s = \epsilon_0 \bar{E}_N$$

$$\bar{E}_N = \frac{\rho_s}{\epsilon_0}$$

According to Gauss law to find the electric flux density at the boundary and electric flux density is zero at the boundary. The surface integral must be calculated for top and lateral.

Blank = $\int_{\text{top}} \bar{D} \cdot \bar{a}_z + \int_{\text{lateral}} \bar{D} \cdot \bar{a}_z + \int_{\text{bottom}} \bar{D} \cdot \bar{a}_z$

$\rho = \int_{\text{top}} \bar{D} \cdot \bar{a}_z$

$\rho = \int_{\text{bottom}} \bar{D} \cdot \bar{a}_z$

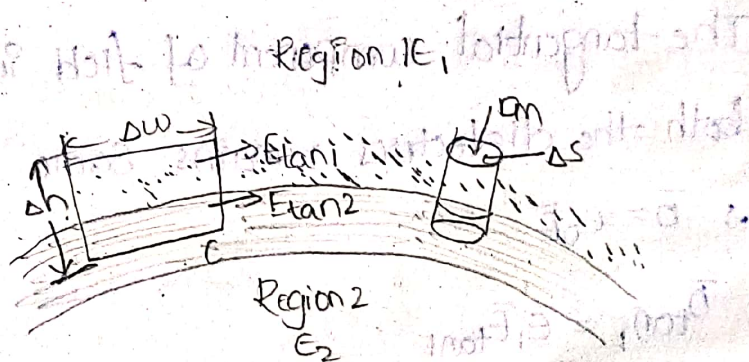
$0 = \int_{\text{lateral}} \bar{D} \cdot \bar{a}_z$

22-10-20
Boundary conditions b/w two dielectrics Two perfect dielectrics

Let us consider (two perfect dielectrics) the boundary between two perfect dielectric. one dielectric has ϵ_1 while the other has permittivity ϵ_2

The \vec{E} and \vec{D} are to obtain again by resolving methods each into two components tangential to the boundary and normal to the surface

Consider a closed path about a rectangular in shape having elementary height Δh & elementary width Δw



$$\oint \vec{E} \cdot d\vec{L} = 0$$

$$\oint \vec{D} \cdot d\vec{S} = Q$$

$$\int_a^b \vec{E} \cdot d\vec{L} + \int_b^c \vec{E} \cdot d\vec{L} + \int_c^d \vec{E} \cdot d\vec{L} + \int_d^a \vec{E} \cdot d\vec{L} = 0$$

$$\therefore \vec{E}_1 = \vec{E}_{1t} + \vec{E}_{1n}$$

$$\vec{E}_2 = \vec{E}_{2t} + \vec{E}_{2n}$$

$$|\vec{E}_{1t}| = E_{tan1} ; |\vec{E}_{2t}| = E_{tan2}$$

$$|\vec{E}_{1n}| = E_{1n} ; |\vec{E}_{2n}| = E_{2n}$$

\int_a^b and \int_c^d becomes zero these are line integrals

along sh and sh $\rightarrow 0$

$$\int_a^b \vec{E} \cdot d\vec{l} + \int_c^d \vec{E} \cdot d\vec{l} = 0$$

$$\vec{E} \int_a^b d\vec{l} + \vec{E} \int_c^d d\vec{l} = 0$$

$$\vec{E} \int_a^b d\vec{l} + \vec{E} \int_c^d d\vec{l} = 0 \quad \therefore \vec{E} \int_a^b d\vec{l} = \vec{E} \int_c^d d\vec{l}$$

$$\vec{E}_{\tan 1} \Delta \omega + \vec{E}_{\tan 2} \Delta \omega = 0$$

$$\vec{E}_{\tan 1} \Delta \omega + \vec{E}_{\tan 2} \Delta \omega = 0$$

$$\vec{E}_{\tan 1} = + \vec{E}_{\tan 2}$$

The tangential component of field intensity in the both the dielectrics remains same

$$\therefore \vec{D} = \epsilon_0 \vec{E}$$

$$\vec{D}_{\tan 1} = \epsilon_1 \vec{E}_{\tan 1}$$

$$\vec{E}_{\tan 1} = \frac{\vec{D}_{\tan 1}}{\epsilon_1}$$

$$\vec{D}_{\tan 2} = \epsilon_2 \vec{E}_{\tan 2}$$

$$\vec{E}_{\tan 2} = \frac{\vec{D}_{\tan 2}}{\epsilon_2}$$

$$\frac{\vec{D}_{\tan 1}}{\vec{D}_{\tan 2}} = \frac{\epsilon_1 \vec{E}_{\tan 1}}{\epsilon_2 \vec{E}_{\tan 2}}$$

$$\frac{\vec{D}_{\tan 1}}{\vec{D}_{\tan 2}} = \frac{\epsilon_1}{\epsilon_2} = \frac{\epsilon_{r1}}{\epsilon_{r2}}$$

The tangential component of \vec{D} undergoes some change across the interface hence tangential \vec{D} is said to be discontinuous across the boundary

To find Normal component

Let us use Gauss law

$$\oint \vec{D} \cdot d\vec{s} = Q$$

Consider a gaussian surface in the form of right circular cylinder

$$\int_{\text{top}} \vec{D} \cdot d\vec{s} + \int_{\text{bottom}} \vec{D} \cdot d\vec{s} + \int_{\text{lateral}} \vec{D} \cdot d\vec{s} = Q$$

$$\therefore \int_{\text{top}} \vec{D} \cdot d\vec{s} + \int_{\text{bottom}} \vec{D} \cdot d\vec{s} = Q \rightarrow \textcircled{1} \quad \because \int_{\text{lateral}} \vec{D} \cdot d\vec{s} = \Delta h = 0$$

$$\int_{\text{top}} \vec{D} \cdot d\vec{s} = \vec{D} \int_{\text{top}} d\vec{s} = \vec{D}_{N1} \Delta S$$

$$\int_{\text{bottom}} \vec{D} \cdot d\vec{s} = \vec{D} \int_{\text{bottom}} d\vec{s} = -\vec{D}_{N2} \Delta S$$

from $\textcircled{1}$

$$\vec{D}_{N1} \Delta S - \vec{D}_{N2} \Delta S = Q$$

$$\therefore \int_S \Delta S = Q$$

$$\vec{D}_{N1} \Delta S - \vec{D}_{N2} \Delta S = \int_S \Delta S$$

$$(\vec{D}_{N1} - \vec{D}_{N2}) \Delta S = \int_S \Delta S$$

$$\rho_s = \vec{D}_{N1} - \vec{D}_{N2}$$

$$\because \rho_s = 0$$

$$\vec{D}_{N1} - \vec{D}_{N2} = 0$$

$$\vec{D}_{N1} = \vec{D}_{N2}$$

$$\therefore \bar{D} = \epsilon \bar{E}$$

$$\bar{D}_{1N} = \epsilon_1 \bar{E}_{1N}$$

$$\bar{D}_{1N} = \epsilon_1 \bar{E}_{1N} \rightarrow \textcircled{1}$$

$$\bar{E}_{1N} = \frac{\bar{D}_{1N}}{\epsilon_1} \rightarrow \textcircled{1}$$

$$\bar{E}_{2N} = \frac{\bar{D}_{2N}}{\epsilon_2} \rightarrow \textcircled{2}$$

$$\bar{D}_{2N} = \epsilon_2 \bar{E}_{2N} \rightarrow \textcircled{2}$$

$$\frac{\textcircled{1}}{\textcircled{2}} = \frac{\bar{D}_{1N}}{\bar{D}_{2N}} = \frac{\epsilon_1 \bar{E}_{1N}}{\epsilon_2 \bar{E}_{2N}}$$

$$\frac{\textcircled{1}}{\textcircled{2}} \frac{\bar{E}_{1N}}{\bar{E}_{2N}} = \frac{\bar{D}_{1N}}{\bar{D}_{2N}} = \frac{\epsilon_1}{\epsilon_2}$$

$$\frac{\bar{E}_{1N}}{\bar{E}_{2N}} = \frac{\epsilon_2}{\epsilon_1}$$

$$\frac{\bar{E}_{1N}}{\bar{E}_{2N}} = \frac{\epsilon_2}{\epsilon_1} = \frac{\epsilon_{r2}}{\epsilon_{r1}}$$

The normal components of dielectric field intensity \bar{E} are inversely proportional to the relative permittivities two media

Hence the normal component of flux density \bar{D} is continuous at the boundary between the two perfect dielectrics

25-10-22
concept of a capacitance:-

Consider two magnetic materials m_1, m_2 which are placed dielectric medium having permittivity ϵ

m_1 carries ~~to~~ positive charge Q . While m_2 carries negative charge equal in magnitude as Q

dielectric field is normal to the conductor surface and the electric flux directed from m_1 towards m_2 in a system.

The potential difference between of m_1 & m_2

The ratio magnitude of the total charge on any one of the two conductors and potential difference between the conductors is called the capacitance of the two conductors. it denoted by C

$$C = \frac{Q}{V}$$

and is measured in faradays (C^2/V) columb/volts

Obtain form. of Gauss law $Q = \oint \vec{D} \cdot d\vec{s}$
using

$$\therefore \vec{D} = \epsilon \cdot \vec{E}$$

$$\therefore Q = \oint \epsilon \cdot \vec{E} \cdot d\vec{s}$$

$$Q = \oint \epsilon \vec{E} \cdot d\vec{s}$$

work done in moving positive charge to negative charge

$$V = - \int_L \vec{E} \cdot d\vec{L}$$

$$V = - \int^+ \vec{E} \cdot d\vec{L}$$

$$C = \frac{q}{\int E \cdot ds}$$

Capacitance in Series:-

$$V = V_1 + V_2 + V_3$$

$$\therefore C = \frac{Q}{V}$$

$$Q = CV$$

$$Q_1 = Q = C_1 V_1 \Rightarrow V_1 = \frac{Q}{C_1}$$

$$Q_2 = Q = C_2 V_2 \Rightarrow V_2 = \frac{Q}{C_2}$$

$$Q_3 = Q = C_3 V_3 \Rightarrow V_3 = \frac{Q}{C_3}$$

$$V = V_1 + V_2 + V_3$$

$$V = \frac{Q}{C_1} + \frac{Q}{C_2} + \frac{Q}{C_3}$$

$$V = Q \left(\frac{1}{C_1} + \frac{1}{C_2} + \frac{1}{C_3} \right)$$

$$\frac{V}{Q} = \frac{1}{C_1} + \frac{1}{C_2} + \frac{1}{C_3} + \dots$$

Capacitance in parallel

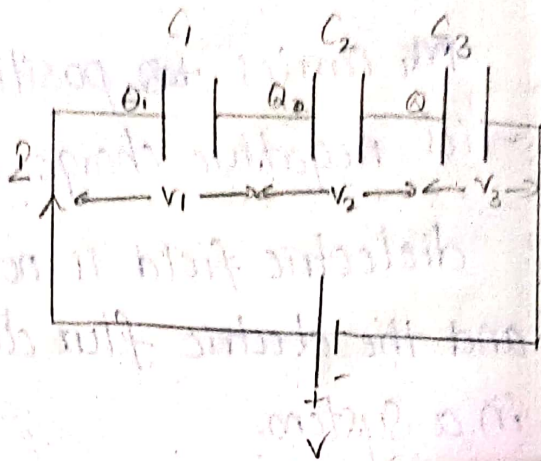
Let us consider the capacitors are connected in parallel C_1, C_2, C_3 across the potential voltage V . In the parallel-connection

the total charge is Q and voltage is same V . The total charge is $Q = Q_1 + Q_2 + Q_3$

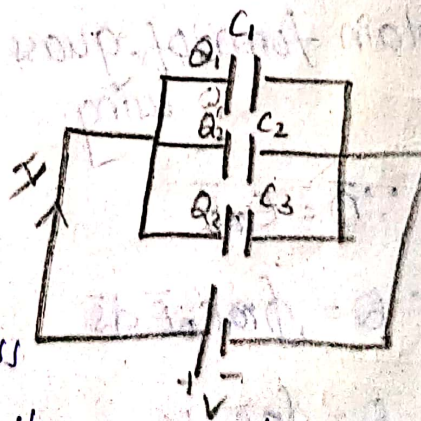
$$\therefore Q = CV$$

$$Q_1 = C_1 V$$

$$Q_2 = C_2 V$$



Let us consider a capacitors are connected in series. the acrossing the voltage V .



$$Q_2 = C_2 V$$

$$Q_3 = C_3 V$$

$$Q = Q_1 + Q_2 + Q_3$$

$$Q = C_1 V + C_2 V + C_3 V$$

$$Q = (C_1 + C_2 + C_3) V$$

$$\frac{Q}{V} = C_1 + C_2 + C_3$$

$$C = C_1 + C_2 + C_3$$

Energy stored in capacitor

Consider a parallel plate capacitor it is supplied with the voltage V .

Let \vec{a}_N is the direction of normal to the plates

$$\vec{E} = \frac{V}{d} \vec{a}_N$$

The energy stored is given by $W_E = \frac{1}{2} \int_{Vol} \vec{D} \cdot \vec{E} dV$

$$W_E = \frac{1}{2} \int_{Vol} \vec{E} \cdot \vec{E} dV$$

$$= \frac{1}{2} \int_{Vol} E \frac{V^2}{d^2} dV$$

$$= \frac{1}{2} \frac{EV^2}{d^2} \int_{Vol} dV$$

$$= \frac{1}{2} \frac{EV^2}{d^2} V$$

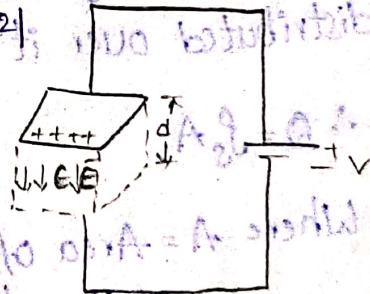
$$= \frac{1}{2} \frac{EV^2}{d^2} (A \times d)$$

$$= \frac{1}{2} \frac{CV^2 A}{d}$$

$$= \frac{1}{2} CV^2$$

$$\vec{E} \cdot \vec{E} = E^2$$

$$E = \frac{V}{d}$$



Energy density

$$W_E = \frac{1}{2} \epsilon \int_{\text{Vol}} |E|^2 dV$$

$$W_E = \frac{1}{2} \epsilon |\bar{E}|^2 \quad \text{J/m}^3$$

$$\bar{D} = \epsilon |\bar{E}| \quad \therefore \bar{E} = \frac{\bar{D}}{\epsilon}$$

$$W_E = \frac{1}{2} \frac{|\bar{D}|^2}{\epsilon} = \frac{1}{2} \bar{D} \cdot \bar{E} \quad \text{J/m}^3$$

$$W_E = \frac{1}{2} \bar{D} \cdot \frac{\bar{D}}{\epsilon}$$

$$W_E = \frac{1}{2} \bar{D} \cdot \bar{E} \quad \text{J/m}^3$$

26-10-22

parallel plate capacitor:

It consists of 2 parallel ^{metallic} plates (capacitor) separated by distance d the space between the plates is filled with a dielectric of permittivity ϵ . The lower plate and upper plate carries with positive charge and negative charge respectively and is distributed over it with a charge density $+\rho_s, -\rho_s$.

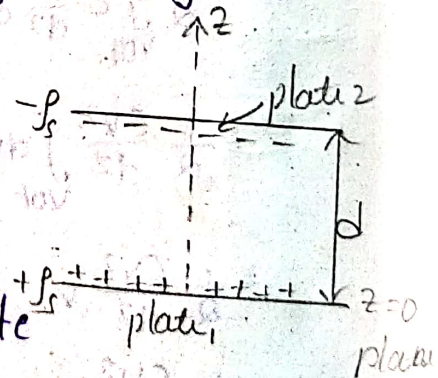
$$\therefore Q = \rho_s A$$

Where A = Area of Cross Section of the plates in m^2

Assuming plate 1 to be infinite sheet of the charge

$$\bar{E}_1 = \frac{\rho_s}{2\epsilon} \bar{a}_N = \frac{\rho_s}{2\epsilon} \bar{a}_z \quad \text{V/m}$$

While for plate 2. $\bar{E}_2 = \frac{-\rho_s}{2\epsilon} \bar{a}_N = \frac{-\rho_s}{2\epsilon} (-\bar{a}_z)$



$$\therefore \vec{E} = \vec{E}_1 + \vec{E}_2$$

$$\vec{E} = \frac{\rho_s}{\epsilon} \vec{a}_z - \frac{\rho_s}{\epsilon} (-\vec{a}_z)$$

$$= \frac{\rho_s}{\epsilon} \vec{a}_z$$

The potential difference is given by $V = -\int_{\text{upper}}^{\text{lower}} \vec{E} \cdot d\vec{L}$

$$\therefore dL = dx \vec{a}_x + dy \vec{a}_y + dz \vec{a}_z$$

$$\therefore V = -\int_{z=d}^{z=0} \frac{\rho_s}{\epsilon} \vec{a}_z \cdot (dx \vec{a}_x + dy \vec{a}_y + dz \vec{a}_z)$$

$$= -\int_d^0 \frac{\rho_s}{\epsilon} dz$$

$$= -\frac{\rho_s}{\epsilon} \int_d^0 dz$$

$$= \frac{\rho_s}{\epsilon} [z]_d^0$$

$$= \frac{\rho_s}{\epsilon} [0 - d]$$

$$= \frac{\rho_s d}{\epsilon}$$

The capacitance of the charge to voltage

$$Q = CV$$

$$C = \frac{Q}{V}$$

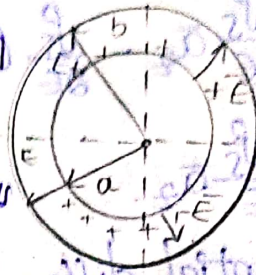
$$C = \frac{\rho_s A}{\frac{\rho_s d}{\epsilon}}$$

$$C = \frac{A\epsilon}{d}$$

$$C = \frac{\epsilon_0 \epsilon_r A}{d}$$

Spherical capacitance:

Consider a spherical capacitor formed two concentric spherical conducting shells of radii a & b



Considering a gaussian surface as a sphere of a radius r it can be obtained by \vec{E} is in radial direction given by $E = \frac{Q}{4\pi\epsilon_0 r^2} \hat{a}_r$ ($\frac{V}{m}$)

The potential difference (betw) is workdone in moving unit positive charge against the direction of \vec{E} .

i.e from b to a .

$$\therefore V = -\int_b^a \vec{E} \cdot d\vec{L} = -\int_b^a E \cdot dr$$

$$= -\int_b^a \frac{Q}{4\pi\epsilon_0 r^2} \hat{a}_r \cdot dr \hat{a}_r$$

$$\therefore d\vec{L} = dr \hat{a}_r$$

$$V = -\int_b^a \frac{Q}{4\pi\epsilon_0 r^2} \hat{a}_r \cdot dr \hat{a}_r$$

$$V = -\frac{Q}{4\pi\epsilon_0} \int_b^a \frac{1}{r^2} dr$$

$$V = -\frac{Q}{4\pi\epsilon_0} \left[\frac{-1}{r} \right]_b^a$$

$$V = -\frac{Q}{4\pi\epsilon_0} \left[\frac{-1}{a} + \frac{1}{b} \right]$$

$$V = \frac{Q}{4\pi\epsilon_0} \left(\frac{1}{a} - \frac{1}{b} \right)$$

$$V = \frac{Q}{4\pi\epsilon} \left(\frac{1}{a} - \frac{1}{b} \right)$$

$$\therefore Q = CV$$

$$C = \frac{Q}{V} = \frac{Q}{\frac{Q}{4\pi\epsilon} \left(\frac{1}{a} - \frac{1}{b} \right)}$$

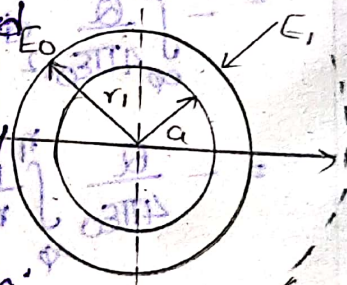
$$C = \frac{Q}{\frac{Q}{4\pi\epsilon} \left(\frac{1}{a} - \frac{1}{b} \right)}$$

$$C = \frac{4\pi\epsilon}{\left(\frac{1}{a} - \frac{1}{b} \right)}$$

14-0-2 Isolated Sphere coated with dielectric

Consider a single isolated coated sphere with a dielectric having permittivity ϵ_1 upto radius r_1 .

The radius of inner sphere is 'a' it is placed in a free space. So outside the sphere ϵ_0 it carries a charge +Q.



$$\therefore E = E_1; a < r < r_1$$

$$E = E_0; r > r_1$$

\therefore potential difference is work done bringing unit positive charge from outer sphere $r = \infty$ to inner sphere $r = a$ against \vec{E}

$$\therefore V = -\int_{\infty}^a \vec{E} \cdot d\vec{L}$$

$$V = -\int_{\infty}^a E \cdot dL$$

$$V = -\int_{\infty}^{r_1} E \cdot dL - \int_{r_1}^a E \cdot dL$$

Now $a < r < r_1$

$$\vec{E}_1 = \frac{Q}{4\pi\epsilon_0 r^2} \vec{a}_r \rightarrow \textcircled{1} \quad \left(\frac{1}{r}\right)' = -\frac{1}{r^2}$$

Now $r > r_1$

$$\vec{E}_2 = \frac{Q}{4\pi\epsilon_0 r^2} \vec{a}_r \rightarrow \textcircled{2} \quad \therefore dL = dr \vec{a}_r$$

$$V = - \int_{\infty}^{r_1} \frac{Q}{4\pi\epsilon_0 r^2} \vec{a}_r \cdot dL - \int_{r_1}^a \frac{Q}{4\pi\epsilon_1 r^2} \vec{a}_r \cdot dL$$

$$= - \int_{\infty}^{r_1} \frac{Q}{4\pi\epsilon_0 r^2} \vec{a}_r \cdot dr \vec{a}_r - \int_{r_1}^a \frac{Q}{4\pi\epsilon_1 r^2} \vec{a}_r \cdot dr \vec{a}_r$$

$$= - \int_{\infty}^{r_1} \frac{Q}{4\pi\epsilon_0 r^2} dr - \int_{r_1}^a \frac{Q}{4\pi\epsilon_1 r^2} dr$$

$$= - \frac{Q}{4\pi\epsilon_0} \int_{\infty}^{r_1} \frac{1}{r^2} dr - \frac{Q}{4\pi\epsilon_1} \int_{r_1}^a \frac{1}{r^2} dr$$

$$= - \frac{Q}{4\pi\epsilon_0} \left[\frac{-1}{r} \right]_{\infty}^{r_1} - \frac{Q}{4\pi\epsilon_1} \left[\frac{-1}{r} \right]_{r_1}^a$$

$$= \frac{+Q}{4\pi\epsilon_0} \left[\frac{1}{r_1} - \frac{1}{\infty} \right] + \frac{Q}{4\pi\epsilon_1} \left[\frac{1}{a} - \frac{1}{r_1} \right]$$

$$= \frac{Q}{4\pi\epsilon_0} \left[\frac{1}{\epsilon_0} \left(\frac{1}{r_1} \right) - 0 \right] + \frac{1}{\epsilon_1} \left(\frac{r_1 - a}{a r_1} \right)$$

$$\frac{Q}{4\pi} \left[\frac{1}{\epsilon_0 r_1} + \frac{1}{\epsilon_1} \left(\frac{r_1 - a}{a r_1} \right) \right]$$

$$= \frac{Q}{4\pi r_1} \left[\frac{1}{\epsilon_0} + \frac{1}{\epsilon_1} \left(\frac{r_1 - a}{a} \right) \right]$$

$$\therefore C = \frac{Q}{V}$$

$$C = \frac{Q}{\frac{Q}{4\pi r_1} \left[\frac{1}{\epsilon_0} + \frac{1}{\epsilon_1} \left(\frac{r_1 - a}{a} \right) \right]}$$

$$C = \frac{4\pi r_1}{\frac{1}{\epsilon_0} + \frac{1}{\epsilon_1} \left(\frac{r_1 - a}{a} \right)}$$

$$\frac{1}{C} = \frac{\frac{1}{\epsilon_0} + \frac{1}{\epsilon_1} \left(\frac{r_1 - a}{a} \right)}{4\pi r_1}$$

$$\frac{1}{C} = \frac{1}{4\pi \epsilon_0 r_1} + \frac{1}{\epsilon_1 4\pi r_1} \left(\frac{1}{a} - \frac{1}{r_1} \right)$$

it compare with $\frac{1}{C} = \frac{1}{C_1} + \frac{1}{C_2}$

$$\frac{1}{C_1} = \frac{1}{4\pi \epsilon_0 r_1}$$

$$C_1 = 4\pi \epsilon_0 r_1$$

$$\frac{1}{C_2} = \frac{\left(\frac{1}{a} - \frac{1}{r_1} \right)}{4\pi \epsilon_1}$$

$$\frac{1}{C} = \frac{1}{C_1} + \frac{1}{C_2}$$

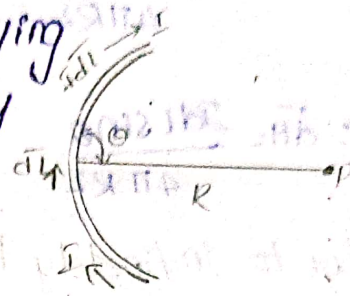
$$C = \frac{1}{\frac{1}{C_1} + \frac{1}{C_2}}$$

$$= \frac{1}{\frac{1}{4\pi \epsilon_0 r_1} + \frac{\left(\frac{1}{a} - \frac{1}{r_1} \right)}{4\pi \epsilon_1}}$$

43 Magneto Statics, Ampere's law and force in Magnetic field

Biot Savart law:

Consider a conductor carrying direct current I , and a study magnetic field produced around it.



- * The biot-Savart law allows to obtain the differential magnetic field intensity dH produced at a point p due to a differential current element $I dL$.
- * It states that the proportional to the product of current (I) and differential length dL .
- * The sine of the angle between the element and the line joining point p to the element and inversely proportional to the square of the distance R between point p and the element.

$$\therefore dH \propto \frac{I dL \sin \theta}{R^2}$$

$$dH = \frac{k I dL \sin \theta}{R^2}$$

where $k = \text{proportionality constant} = \frac{1}{4\pi}$

$$dH = \frac{1}{4\pi} \frac{I dL \sin \theta}{R^2}$$

$\vec{a}_r =$ unit vector in the direction from differential current element to point p then from cross

product vector

$$d\vec{L} \times \vec{a}_{R_2} = |dL| |\vec{a}_{R_2}| \sin \theta$$

$$= (dL) \sin \theta$$

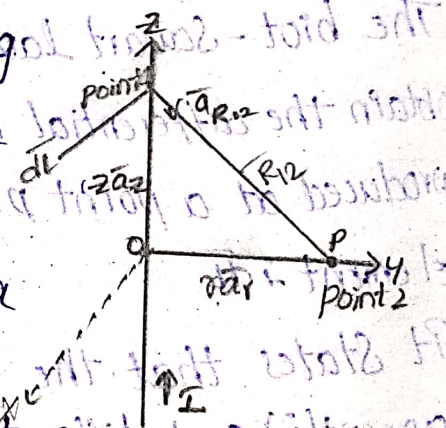
$$\therefore dH = \frac{I dL \times \vec{a}_{R_2}}{4\pi R^2}$$

$$\therefore dH = \frac{I dL \sin \theta}{4\pi R^2}$$

7281
x 281

H due to infinitely long straight conductor

Consider an infinitely long straight conductor along z-axis. The current passing through the conductor is a direct current I amp. The field intensity H at a point P which is at a distance r from the z-axis. Small differential element at a point 1 along the z-axis at a distance z from the origin.

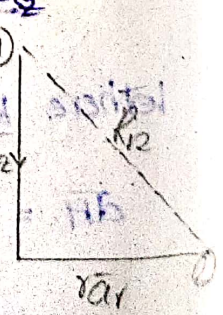


$$\therefore d\vec{L} = dz \vec{a}_z$$

The distance vector joining point 1 to point 2 is \vec{R}_{12} & can be written as $\vec{R}_{12} = r\vec{a}_r - z\vec{a}_z = \vec{R}$

$$\vec{a}_{R_{12}} = \frac{(\vec{R}_{12})}{|\vec{R}_{12}|} = \frac{r\vec{a}_r - z\vec{a}_z}{\sqrt{r^2 + z^2}}$$

$$d\vec{L} \times \vec{a}_{R_{12}} = \begin{vmatrix} \vec{a}_1 & \vec{a}_2 & \vec{a}_3 \\ 0 & 0 & dz \\ 0 & 0 & -z \end{vmatrix}$$



1-1/2

$$d\vec{l} \times \vec{a}_{R_{12}} = -a_{\phi}(0 - r dz)$$

$$d\vec{l} \times \vec{a}_{R_{12}} = r dz a_{\phi}$$

$$\therefore d\vec{l} \times \vec{a}_{R_{12}} = \frac{r dz a_{\phi}}{\sqrt{r^2 + z^2}}$$

According to biot savart's law $d\vec{H}$ at point 2 is

$$d\vec{H} = \frac{2 d\vec{l} \times \vec{a}_{R_{12}}}{4\pi R_{12}^2} \text{ (amp-m)}$$

$$\therefore d\vec{H} = \frac{2 r dz a_{\phi}}{\sqrt{r^2 + z^2} 4\pi R_{12}^2}$$

$$d\vec{H} = \frac{2 r dz a_{\phi}}{4\pi \sqrt{r^2 + z^2} (\sqrt{r^2 + z^2})^2}$$

$$d\vec{H} = \frac{2 r dz a_{\phi}}{4\pi (r^2 + z^2)^{3/2}} \rightarrow \textcircled{1}$$

$$\therefore \text{put } z = r \tan \theta \Rightarrow z^2 = r^2 \tan^2 \theta$$

$$dz = r \sec^2 \theta d\theta \quad \theta \text{ limits: } -\frac{\pi}{2} \text{ to } \frac{\pi}{2}$$

Integrating eq ①

$$H = \int_{-\pi/2}^{\pi/2} \frac{2 r dz a_{\phi}}{4\pi (r^2 + z^2)^{3/2}}$$

$$H = \int_{-\pi/2}^{\pi/2} \frac{2 r r^2 \sec^2 \theta d\theta a_{\phi}}{4\pi (r^2 + r^2 \tan^2 \theta)^{3/2}}$$

$$H = a_{\phi} \int_{-\pi/2}^{\pi/2} \frac{2 r^3 \sec^2 \theta d\theta}{4\pi (r^2 + r^2 \tan^2 \theta)^{3/2}}$$

$$\vec{H} = a_\phi \int_{-\pi/2}^{\pi/2} \frac{I \sec^2 \theta d\theta}{4\pi r^3 (1 + \tan^2 \theta)^{3/2}}$$

$$\vec{H} = a_\phi \int_{-\pi/2}^{\pi/2} \frac{I \sec^2 \theta d\theta}{4\pi r \sec^3 \theta}$$

$$\vec{H} = a_\phi \int_{-\pi/2}^{\pi/2} \frac{I}{4\pi r} \cos \theta d\theta$$

$$\vec{H} = a_\phi \left[\frac{I}{4\pi r} \sin \theta \right]_{-\pi/2}^{\pi/2}$$

$$\vec{H} = a_\phi \frac{I}{4\pi r} [\sin \pi/2 - \sin(-\pi/2)]$$

$$\vec{H} = a_\phi \frac{I}{4\pi r} (1 - (-1))$$

$$\vec{H} = a_\phi \frac{I}{2\pi r}$$

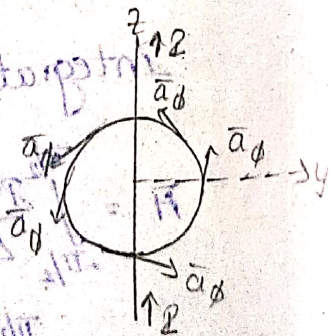
$$\vec{H} = \frac{I}{2\pi r} a_\phi$$

Ampere's circuit law:-

In magnetostatics, the complex problems can be solved using law called ampere's circuit law (or) Ampere's work law

it states that the line integral of magnetic field intensity \vec{H} around along a closed path is exactly equal to the direct current enclosed by that path

$$\therefore \oint \vec{H} \cdot d\vec{L} = I$$



proof of Ampere circuit law:

Consider a long straight conductor carrying direct current I placed along z axis. Radius ' r ', The point P is at a perpendicular distance r from the conductor. Consider dL at a point P which is in \bar{a}_ϕ direction, Tangential to circular path at point P .

$$\therefore dL = r d\phi \bar{a}_\phi$$

\therefore Biot Savart law due to infinitely long conductor

$$H = \frac{I}{2\pi r} \bar{a}_\phi$$

$$H \cdot dL = \frac{I}{2\pi r} \bar{a}_\phi \cdot r d\phi \bar{a}_\phi = \frac{I}{2\pi} r d\phi \quad \therefore \bar{a}_\phi \cdot \bar{a}_\phi = 1$$

$$= \frac{I}{2\pi} d\phi$$

$$\oint H \cdot dL = \int_0^{2\pi} \frac{I}{2\pi} d\phi$$

$$= \frac{I}{2\pi} [\phi]_0^{2\pi}$$

$$= \frac{I}{2\pi} (2\pi - 0)$$

$$\oint H \cdot dL = I$$

Applications of Amperes work law:

(1) H due to infinitely long straight conductor.

$$\therefore H = H \bar{a}_\phi$$

$$dL = r d\phi \bar{a}_\phi$$

$$H \cdot dL = H \bar{a}_\phi \cdot r d\phi \bar{a}_\phi$$

$$H \cdot dL = H r d\phi$$

According to ampere work law $\oint H \cdot dL = I$

$$\int_0^{2\pi} H_\phi r d\phi = I$$

$$H_\phi r [0]_0^{2\pi} = I$$

$$H_\phi r (2\pi - 0) = I$$

$$H_\phi = \frac{I}{2\pi r}$$

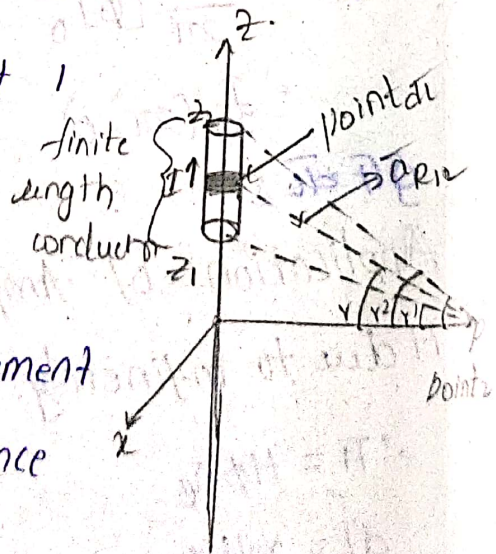
Hence H_ϕ at a point p is given by

$$H = H_\phi \bar{a}_\phi$$

$$H = \frac{I}{2\pi R} \bar{a}_\phi \text{ A/m}$$

Magnetic field intensity due to straight conductor of finite length

Consider a conductor of finite length placed along z -axis. it carries a direct current "i", the \perp distance of point p from z -axis is "r" the conductor is placed such that 1 and is at $z=z_1$, while the other is at $z=z_2$

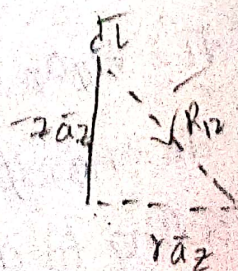


Consider a differential element dl along z -axis at a distance z from origin therefore

$$d\vec{L} = dz \bar{a}_z$$

$$\bar{a}_{R12} = \frac{R_{12}}{|R_{12}|}$$

$$\bar{R}_{12} = r\bar{a}_r - z\bar{a}_z$$



$$S_0 = \frac{r a R - z a_2}{\sqrt{r^2 - z^2}}$$

$$dF \times \bar{a}_{R12} = r d z a \phi$$

by det

$$\int dL \times \bar{a}_{R12} = \frac{\int r d z a \phi}{\sqrt{r^2 + z^2}}$$

$$\begin{vmatrix} r & \phi & z \\ 0 & 0 & dz \\ r & 0 & -z \end{vmatrix}$$

According to Biot Savart's law $a \phi (0-r dz)$

$$dH = \frac{\int dL \times \bar{a}_{R12}}{4\pi R_{12}^2}$$

$$dH = \frac{\int r d z a \phi}{4\pi \sqrt{r^2 + z^2} (\sqrt{r^2 + z^2})^2}$$

$$= \frac{\int r d z a \phi}{4\pi (r^2 + z^2)^{3/2}}$$

put $z = r \tan \alpha$ with limits α_1 to α_2

$$z^2 = r^2 \tan^2 \alpha$$

$$z = z_1 = r \tan \alpha_1$$

$$dz = r \sec^2 \alpha d\alpha$$

$$z = z_2 = r \tan \alpha_2$$

Applying integration for above eq

$$\int dH = H = \int_{\alpha_1}^{\alpha_2} \frac{\int r d z r^2 \sec^2 \alpha d\alpha a \phi}{4\pi (r^2 + r^2 \tan^2 \alpha)^{3/2}}$$

$$= \int_{\alpha_1}^{\alpha_2} \frac{\int r^2 \sec^2 \alpha d\alpha a \phi}{4\pi (r^2)^{3/2} (1 + \tan^2 \alpha)^{3/2}}$$

$$= \int_{\alpha_1}^{\alpha_2} \frac{\int r \sec^2 \alpha d\alpha a \phi}{4\pi (\sec^2 \alpha)^{3/2}}$$

$$= \int_{\alpha_1}^{\alpha_2} \frac{\int r d\alpha a \phi}{4\pi r \sec \alpha}$$

$$= \frac{I}{4\pi r} a_\phi \int_{\alpha_1}^{\alpha_2} \cos \alpha d\alpha$$

$$= \frac{I}{4\pi r} a_\phi (\sin \alpha)_{\alpha_1}^{\alpha_2}$$

$$H = \frac{I}{4\pi r} (\sin \alpha_2 - \sin \alpha_1) a_\phi$$

formula for magnetic flux density

$$B = \mu H = \frac{\mu I}{4\pi r} (\sin \alpha_2 - \sin \alpha_1) a_\phi \left(\frac{wb}{m} \right)$$

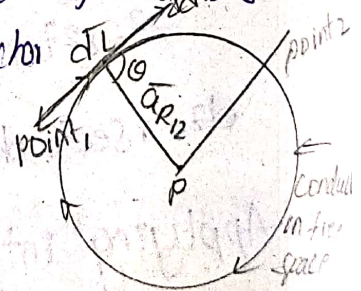
Magnetic field intensity at the centre of circular conductor

Consider a current carrying conductor arranged in a circular form. The H at the centre of the circular loop is to be obtained.

Consider a differential length dL at a point P is tangential to the circular conductor at point 1 where

$\theta =$ angle b/w dL & \bar{a}_{R2}

- Cross product $I dL \times \bar{a}_{R2} = I dL |a_{R2}| \sin \theta$
 $= I dL \sin \theta \bar{a}_N$



According to biot savarts law

$$dH = \frac{I dL \times \bar{a}_{R2}}{4\pi R^2} = \frac{I dL \sin \theta \bar{a}_N}{4\pi R^2}$$

$$\oint dH = H = \frac{I \sin \theta}{4\pi R^2} \bar{a}_N \oint dL$$

where $\oint dL = 2\pi R$ is the circumference of the circle

Replacing $d\vec{l}$ with $2\pi R$

$$= \frac{\lambda \sin\theta}{4\pi R^2} \bar{a}_N 2\pi R$$

$$\vec{H} = \frac{\lambda \sin\theta}{2R} \bar{a}_N$$

As $d\vec{l}$ is the tangential to the circle R_1 or R_2 radius, Angle θ , $\angle O$ must be 90°

$$\vec{H} = \frac{\lambda \sin 90^\circ}{2R} \bar{a}_N$$

$$\vec{H} = \frac{\lambda}{2R} \bar{a}_N \left(\frac{A}{m}\right)$$

Magnetic flux density

$$B = \mu H = \frac{\mu_0 I}{2R} \bar{a}_N \text{ Wb/m}^2$$

Magnetic field intensity due to infinite sheet of current:-

- * Consider an infinite sheet of current in the $z=0$ plane, the surface current density is K . The current is flowing in y -direction, hence $K = Ky\bar{a}_y$
- * Consider a closed path 1, 2, 3, 4 the width of the path is 'b' the height is 'a'
- * It is perpendicular to the direction of currents hence in xz plane
- * The current flowing across the distance 'b' is given by
 $\therefore I_{\text{enclosed}} = Kyb$
- * Consider the magnetic lines of force due to the current in \bar{a}_y direction according to right hand thumb rule
- * As current is flowing in y -direction, it cannot have component in y -direction. So, \vec{H} as only

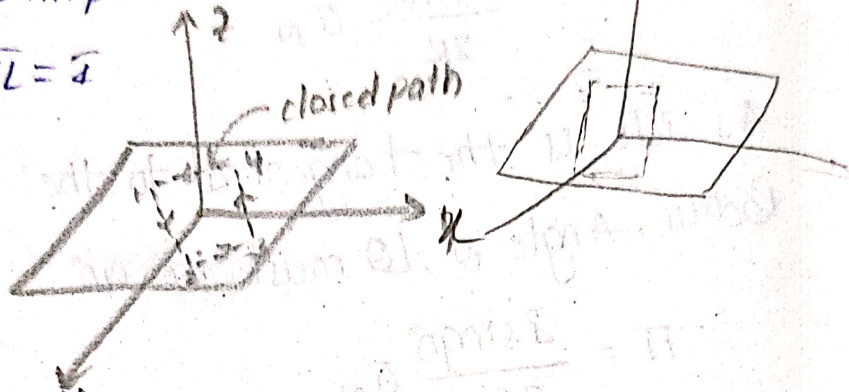
component in x-direction

$$H = H_x \bar{a}_x \text{ for } z > 0$$

$$H = -H_x \bar{a}_x \text{ for } z < 0$$

According to Amphere circuit law

$$\oint \bar{H} \cdot d\bar{L} = I$$



Evaluate the integral along the path 1-2-3-4-1
for path 1-2 $d\bar{L} = dz \bar{a}_z$ since 3-4 is same. \bar{H} is in x direct
ion by while $\bar{a}_x \cdot \bar{a}_z = 0$. Hence 1-2 & 3-4 the integral

$$\oint \bar{H} \cdot d\bar{L} = 0$$

Consider path 2-3 along which $d\bar{L} = dx \bar{a}_x$

$$\begin{aligned} \int_2^3 \bar{H} \cdot d\bar{L} &= \int_2^3 (-H_x \bar{a}_x) (dx \bar{a}_x) \\ &= -H_x \int_2^3 dx = -H_x b \end{aligned}$$

\therefore 1-4, are 4-1

$$\int_4^1 \bar{H} \cdot d\bar{L} = b H_x \quad \therefore \oint \bar{H} \cdot d\bar{L} = b H_x + b H_x = 2b H_x$$

$$\therefore I_{enc} = kyb$$

$$2b H_x = kyb$$

$$H_x = \frac{kyb}{2b} \quad H_x = \frac{1}{2} ky$$

An infinite sheet of current density K (A/m)

$$\bar{H} = \frac{1}{2} K \times \bar{a}_N$$

where \bar{a}_N = unit vector normal from from the current sheet to the point at which \bar{H} is to be obtained

Force on a moving point charge: (Lorentz force eqn)

for a positive charge the force exerted on it is in the direction of \bar{E} . This force is also called as electric force

A point charge of Q

$Q = -1.2 \mu\text{C}$ has a velocity

$\vec{v} = 5\vec{a}_x + 2\vec{a}_y + 3\vec{a}_z$ m/s. Find the magnitude of

the force, the charge if $\vec{E} = (-18\vec{a}_x + 5\vec{a}_y - 10\vec{a}_z) \frac{\text{V}}{\text{m}}$

$\vec{B} = (-4\vec{a}_x + 4\vec{a}_y + 3\vec{a}_z)$ Tesla. Both are present in

simultaneously

$$F_e = Q\vec{E}$$

$$= (-1.2)(-18\vec{a}_x + 5\vec{a}_y - 10\vec{a}_z)$$

$$= 21.6\vec{a}_x - 6\vec{a}_y + 12\vec{a}_z$$

$$F_m = Q(\vec{v} \times \vec{B})$$

$$= (-1.2)(5\vec{a}_x + 2\vec{a}_y + 3\vec{a}_z \times (-4\vec{a}_x + 4\vec{a}_y + 3\vec{a}_z))$$

$$= -1.2(-20\vec{a}_x - 24\vec{a}_y + 36\vec{a}_z)$$

$$= (24\vec{a}_x + 28.8\vec{a}_y - 43.2\vec{a}_z)$$

$$= 24\vec{a}_x + 28.8\vec{a}_y - 43.2\vec{a}_z$$

$$\begin{vmatrix} \vec{a}_x & \vec{a}_y & \vec{a}_z \\ -6 & -24 & 36 \\ -4 & 4 & 3 \end{vmatrix}$$

$$= \vec{a}_x(-7-2-14 \cdot 4) - \vec{a}_y(-18+14 \cdot 4) + \vec{a}_z(-24-96)$$

$$= \vec{a}_x(-21.6) - \vec{a}_y(-13.6) + \vec{a}_z(-33.6)$$

Magnitude of electric force is given by $|F_e|$

$$|F_e| = \sqrt{(21.6)^2 + (-6)^2 + (12)^2}$$

$$|F_e| = \sqrt{466.56 + 36 + 144} = 25.42$$

Magnitude of magnetic force is given by

$$|F_m| = \sqrt{(21.6)^2 + (3.6)^2 + (-33.6)^2}$$

$$|F_m| = 40.101 \text{ N}$$

$$\vec{F}_0 = \vec{F}_e + \vec{F}_m$$

$$x = 25.42$$

$$x =$$

$$\vec{a} \times \vec{b} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ a_x & a_y & a_z \\ b_x & b_y & b_z \end{vmatrix}$$

$$\vec{F} = \vec{F}_e + \vec{F}_m$$

$$\vec{F} = 21.6\hat{a}_x - 6\hat{a}_y + 12\hat{a}_z + -21.6\hat{a}_x - 3.6\hat{a}_y - 33.6\hat{a}_z$$

$$\vec{F} = -2.4\hat{a}_y - 21.6\hat{a}_z$$

$$|F| = \sqrt{(-2.4)^2 + (21.6)^2}$$

$$|F| = 21.70$$

Force on a differential current element

The force exerted on a differential element of charge dq moving in steady magnetic field is given by $d\vec{F} = dq\vec{v} \times \vec{B}$ (N) \rightarrow ①

$$\text{Since } \vec{J} = \rho_v \vec{v} \rightarrow$$
 ②

Current density

$$\text{differential element of charge } dq = \rho_v dV \rightarrow$$
 ③

Sub ③ in ① we get

$$d\vec{F} = \rho_v dV \vec{v} \times \vec{B}$$

$$d\vec{F} = \vec{J} \times \vec{B} dV$$

The force exerted on surface current density is given by $d\vec{F} = \vec{K} \times \vec{B} dS$

Similarly the exerted on differential current Element is given by $-I \vec{L} \times \vec{B} = I d\vec{L}$

$$d\vec{F} = I d\vec{L} \times \vec{B}$$

integrate above eq over a closed path we get

$$\boxed{\vec{F} = \oint I d\vec{L} \times \vec{B}}$$

Force on a straight long current conductor:

If a conductor is straight & long carrying current conductor and the field \vec{B} is uniform along it, then integrating differential force $d\vec{F}$ represented in above equation we get simple expression for the force as

$$\vec{F} = I \vec{L} \times \vec{B}$$

The Magnitude of the force is given by

$$F = ILB \sin \theta$$

The magnetic field exerts a magnetic force on the electrons which constitute the current 'I'

Prob: A conductor 6m long, lies along z-direction with a current of 2Amp in $+\hat{z}$ direction. Find the force experienced by conductor if $\vec{B} = 0.08 \hat{x}$ Tesla.

$$\vec{F} = I d\vec{L} \times \vec{B}$$

$$F = 2(6)\hat{a}_z \times 0.08\hat{a}_x$$

$$F = 12 \times 0.08 \hat{a}_y$$

08-11-22

$F = 0.96 \text{ ay (Newtons)}$

0.08
1.2
0.16
0.8

Force between different current element:

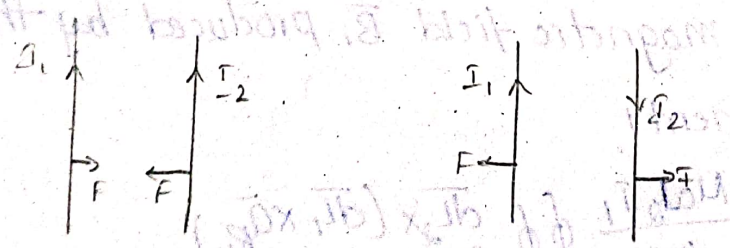
Consider that the two current carrying conductors are placed parallel to each other. Each of this conductor produces its own flux around it. When such two conductors placed closed to each other, there exists a force due to the interaction of two fluxes

The force between the such parallel I carrying conductors depends on the direction of both currents

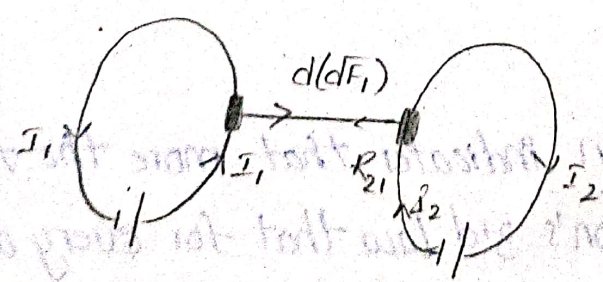
If the direction of both currents are same then the conductors experience a force of attraction

If the direction of two currents are opposite to each other then the conductors experience a force of repulsion

Consider two current elements $I_1 dL_1$ & $I_2 dL_2$



Force between two parallel current carrying conductor



Force between two current elements

The force exerted on a differential current element is given by

$$d(\overline{dF}_1) = I_1 d\overline{L}_1 \times d\overline{B}_2$$

$$d\overline{B}_2 = \mu_0 I_2 \overline{dL}_2 = \mu_0 \left[\frac{I_2 d\overline{L}_2 \times \overline{a}_{R_{21}}}{4\pi R_{21}^2} \right]$$

$$d(\overline{dF}_1) = I_1 d\overline{L}_1 \times \mu_0 \left[\frac{I_2 d\overline{L}_2 \times \overline{a}_{R_{21}}}{4\pi R_{21}^2} \right] \rightarrow \text{---}$$

eqn (1) integrating twice the total force \overline{F}_1 on current element 1 and due to current element 2 is given by

$$\overline{F}_1 = \frac{\mu_0 I_1 I_2}{4\pi} \iint_{L_1 L_2} \frac{d\overline{L}_1 \times (d\overline{L}_2 \times \overline{a}_{R_{21}})}{R_{21}^2}$$

$$\overline{F}_2 = \frac{\mu_0 I_2 I_1}{4\pi} \iint_{L_1 L_2} \frac{d\overline{L}_2 \times (d\overline{L}_1 \times \overline{a}_{R_{12}})}{R_{12}^2}$$

Similarly \overline{F}_2 exerted on the current element 2 due to the magnetic field \overline{B}_1 produced by the current element 1.

$$\overline{F}_2 = \frac{\mu_0 I_2 I_1}{4\pi} \iint_{L_2 L_1} \frac{d\overline{L}_2 \times (d\overline{L}_1 \times \overline{a}_{R_{12}})}{R_{12}^2}$$

$$\therefore \overline{F}_2 = -\overline{F}_1$$

The above condition indicates that more the forces \overline{F}_1 & \overline{F}_2 obey Newton's 3rd law that for every action there is equal & opposite reaction.

Force between two straight, long straight & parallel conductors carrying currents

consider two parallel long straight conductors of length l each carrying current I_1 & I_2

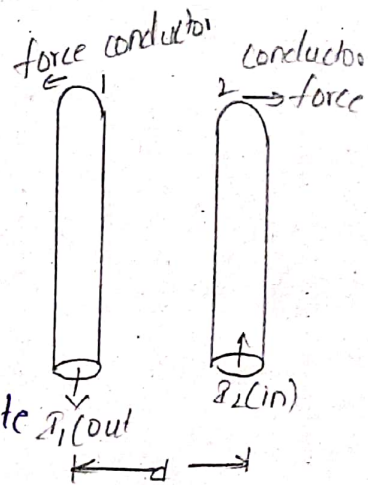
let d be the distance of separation between the two conductors. the current in conductor 1 is moving out while that through conductor 2 is moving in. thus the two currents are in opposite direction

If the directions of the currents through the conductors attract & are same. then the two conductors attract each other while if the direction of currents through the conductors are opposite

The two conductors repel each other

(c) The force exerted on a conductor is given by

$$F = \frac{\mu_0 I_1 I_2 l}{2\pi d} \text{ Newtons}$$



Self and Mutual Inductance

Inductance:-

It is the property of the Material which opposes the rate of change of current to pass through it.

Self inductance:-

The property of a coil that opposes the change in current through it is called self inductance. It is denoted by the letter 'L' and is measured in Henry's.

x] When a closed path or circuit carries a current (I), a magnetic field (B) is produced. This causes a magnetic flux (φ) which is given by $\phi = \oint \vec{B} \cdot d\vec{s}$

x] The flux linkage is defined as the product no. of turns and the total flux linking each of the turns. It is denoted by λ and is measured in weber turn.

$$\therefore \lambda = N\phi \text{ (w-turn)}$$

x] The ratio of the total flux linkages to the current flowing through the circuit is called inductance and is given by

$$L = \frac{N\phi}{I} \text{ (weber turn/Amp)} \rightarrow \textcircled{1}$$

x] According to self induced emf

$$e = -L \frac{di}{dt}$$

$$L = \frac{e}{\left(\frac{di}{dt}\right)} \rightarrow \textcircled{2}$$

Method - 10:

$$\phi = \frac{\text{mmf}}{\text{reluctance}}$$

$$\phi = \frac{NI}{\frac{l}{\mu_0 \mu_r a}}$$

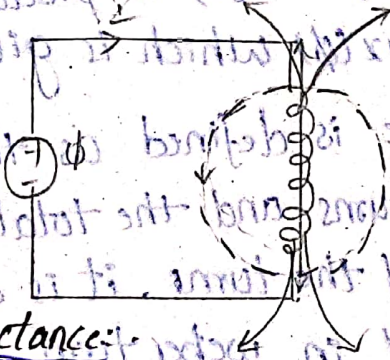
$$\phi = \frac{NI \mu_0 \mu_r a}{l} \rightarrow (3)$$

Now Eq (3) in eq (1)

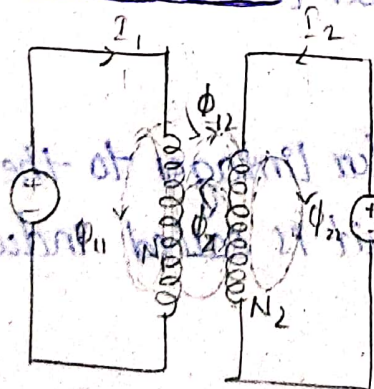
$$L = \frac{N}{I} \left(\frac{NI \mu_0 \mu_r a}{l} \right)$$

$$L = \frac{N^2 \mu_0 \mu_r a}{l}$$

$$L = \frac{N^2}{\int \frac{dl}{\mu}}$$



Mutual inductance:



- *] The flux produced by circuit 1 due to current I_1 flowing through it is denoted by ϕ_{11} & ϵ_1
- *] Similarly the flux produced by circuit 2 due to current I_2 flowing through it is denoted by ϕ_{22}

→ The flux produced by links with the circuit itself and the other circuit. So the flux ϕ_{11} that links with the circuit 2 it denoted by ϕ_{12} and the flux ϕ_{22} that links with the circuit 1 it denoted by ϕ_{21} .

→ The Mutual inductance b/w the two circuits is defined as the flux linkages of one circuit to the current in other circuit.

$$M_{12} = \frac{N_2 \phi_{12}}{I_1}$$

$$M_{21} = \frac{N_1 \phi_{21}}{I_2} \quad (H)$$

$$E_m = M \frac{dI_1}{dt}$$

$$M = \frac{E_m}{\left(\frac{dI_1}{dt}\right)} \quad (H)$$

$$M = \frac{N_2 \phi_{12}}{I_1}$$

$$= \frac{N_2}{\mu_1} \frac{N_1 \mu_1 \mu_0 N_1 r a}{2l}$$

$$= \frac{N_1 N_2 \mu_0 N_1 r a}{2l}$$

$$= \frac{N_1 N_2}{S} \quad (H)$$

02-11-22

2 Coefficient of coupling b/w two circuits

When the two magnetic circuits kepts closed to each interacts with each magnetically through the flux linkages in the circuit due to the current in other circuit, then the circuits are called Magnetically couple circuits.

Self inductance of coil 1 $L_1 = \frac{N_1 \phi_1}{I_1}$

self inductance of coil 1 = $\frac{N_1 \phi_1}{I_1} \rightarrow (2)$

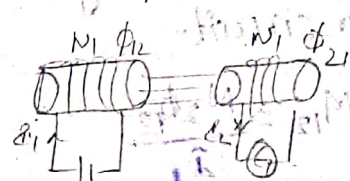
Mutual inductance coil 1 = $\frac{N_2 \phi_{12}}{I_1} \rightarrow (3)$

Mutual inductance coil 2 = $\frac{N_1 \phi_{21}}{I_2} \rightarrow (4)$

$\phi_{12} = k_1 \phi_1$ $\phi_{21} = k_2 \phi_2$

(3) $\Rightarrow M_{12} = \frac{N_2 k_1 \phi_1}{I_1} \rightarrow (5)$

(4) $\Rightarrow M_{21} = \frac{N_1 k_2 \phi_2}{I_2} \rightarrow (6)$



eq (5) x eq (6)

$M_{12} \times M_{21} = \frac{N_2 k_1 \phi_1}{I_1} \times \frac{N_1 k_2 \phi_2}{I_2}$

$M \times M = \frac{N_1 N_2 k_1 k_2 \phi_1 \phi_2}{I_1 I_2}$

$M^2 = k_1 k_2 \left(\frac{N_1 \phi_1}{I_1} \right) \left(\frac{N_2 \phi_2}{I_2} \right)$

$M^2 = k_1 k_2 L_1 L_2$

$M^2 = k^2 L_1 L_2$

$k^2 = \frac{M^2}{L_1 L_2}$

$k = \sqrt{\frac{M^2}{L_1 L_2}}$

$k = \frac{M}{\sqrt{L_1 L_2}}$

*] The two magnetic circuits are coupled together in series Aided and if m is mutual inductance between them then the effective inductance of the system is given by

$L = L_1 + L_2 + 2M$ Henry

Similarly two magnetic circuits L_1 & L_2 are magnetically coupled in series opposing then the effective inductance of the series system is given by

$$L_{eq} = L_1 + L_2 - 2M$$

parallel Aiding $L_{eq} = \frac{L_1 L_2 - M^2}{L_1 + L_2 - 2M}$ *Mutually repulsive*

parallel opposing $L_{eq} = \frac{L_1 L_2 - M^2}{L_1 + L_2 + 2M}$

06-12-22

Pb: A Solenoid with $n_1 = 2000$, $r_1 = 2\text{ cm}$ & $L_1 = 100\text{ cm}$ is concentric with in a second coil of $n_2 = 4000$, $r_2 = 4\text{ cm}$; $L_2 = 100\text{ cm}$. Find mutual inductance assuming free space inductance

Given $N_1 = 2000$

$r_1 = 2\text{ cm}$

$L_1 = 100\text{ cm}$

$N_2 = 4000$

$r_2 = 4\text{ cm}$

$L_2 = 100\text{ cm}$

$$M = \frac{N_1 N_2}{\frac{l}{\mu_0 \mu_r \mu_a}}$$

$$M = \frac{N_1 N_2 \mu_0 \mu_r \mu_a}{l}$$

$$a = \frac{\pi d^2}{4}$$

$$d = 2r$$

Magnetic field intensity $H_1 = \frac{N_1 I_1}{L_1}$

$$= \frac{2000 \times 2}{100}$$

$$= 20 \text{ A/m} = 20 \times 10^{-2} \text{ A/m} = 2000 \text{ A/m}$$

Magnetic field density $B = \mu H$

$B = \mu H$

$B_1 = \mu H_1$

$B_1 = \mu_0 \mu_r \mu_a H_1 = 4\pi \times 10^{-7} \times 1 \times 2000 \text{ A/m}$



Total flux $\phi_1 = B_1 A$

$$\phi_1 = 2.512 \times 10^{-3} B_1 \times [\pi (2 \times 10^{-2})^2]$$

$$\phi_1 = 2.512 \times 10^{-3} \times \pi^2 \times 4 \times 10^{-4}$$

$$\phi_1 = 3.15 \times 10^{-6} I_1 \text{ weber}$$

$$M_{21} = \frac{N_2 \phi_{12}}{I_1} = \frac{400 \times 3.15 \times 10^{-6} I_1}{I_1} = 0.0126 \text{ H}$$

pb: If a coil of $800 \mu\text{H}$ is magnetically couple to another coil of $200 \mu\text{H}$. The co-efficient of coupling betw^o coils is 0.05. calculate inductance of two coils are connect-
 -ed (a) Series Aided (b) Series opposing (c) parallel aided
 (d) parallel opposing

Given $L_1 = 800 \times 10^{-6} \text{ H}$

$L_2 = 200 \times 10^{-6} \text{ H}$

$k = 0.05$

(a) Series aided

$$L_{eq} = L_1 + L_2 + 2M$$

$$k = \frac{M}{\sqrt{L_1 L_2}}$$

$$M = k \sqrt{L_1 L_2}$$

$$M = 2 \times 10^{-5} \text{ H}$$

$$L_{eq} = 800 \times 10^{-6} + 200 \times 10^{-6} + 2 \times 2 \times 10^{-5} =$$

$$= 1.04 \times 10^{-3} \text{ H}$$

(b) Series opposing

$$L_{eq} = L_1 + L_2 - 2M$$

$$L_{eq} = 1000 \times 10^{-6} - 2 \times 2 \times 10^{-5} = 9.6 \times 10^{-4} \text{ H}$$

(c) parallel aidnid

$$L_{eq} = \frac{L_1 L_2 - M^2}{L_1 + L_2 - 2M}$$

$$= \frac{800 \times 10^{-6} \times 200 \times 10^{-6} - (2 \times 10^{-5})^2}{9.6 \times 10^{-4}}$$

120
100
200
100
300
56
350
400

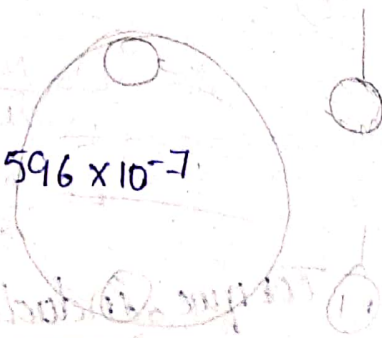
~~1.596 x 10^-7~~
~~1.53 x 10^-4~~
 = 1.6625 x 10^-2 H

(d) parallel opposing

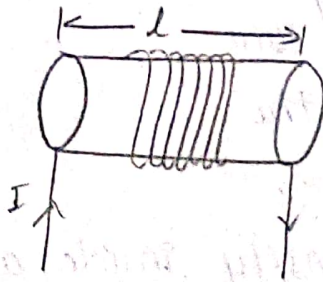
$$L_{eq} = \frac{L_1 L_2 - M^2}{L_1 + L_2 + 2M}$$

$$= 1.596 \times 10^{-7}$$

$$= 1.53 \times 10^{-4} \text{ H}$$



Inductance of a solenoid:



$\frac{NI}{l}$
 $B = \mu_0 H$
 $\Phi = \frac{NI}{l} l$
 $\Phi = N\phi$
 $\phi = \frac{\Phi}{A}$
 $\Phi = NBA$

Consider a solenoid of (N) turns, let the current flowing through the solenoid be I. let the length of the solenoid 'l' & cross section Area be A (m²)

*] The field intensity inside of the solenoid is given by $H = \frac{NI}{l}$ ($\frac{A \text{ turns}}{m}$)

* Total flux linkages = $N\Phi$
 $= NBA$
 $= NMHA$
 $= NNA \left(\frac{NI}{l} \right)$

$B = \frac{\Phi}{A}$
 $B = \mu_0 H$

$$\text{Total flux linkages} = \frac{\mu N^2 A I}{l}$$

The inductance of a solenoid is given by

$$L = \frac{\text{Total flux linkages}}{\text{Total current}}$$

$$L = \frac{\mu N^2 A I}{l}$$

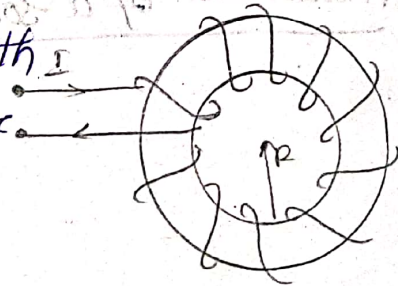
$$L = \frac{\mu N^2 A}{l} \text{ Henry}$$

$$L = \frac{N^2}{\frac{l}{\mu A}}$$

$$L = \frac{N^2}{S} \text{ Henry}$$

Torque Inductance of Toroid:-

Consider a toroidal ring with N turns and carrying current I . Let the radius of the toroid be R .



The magnetic flux density inside a toroidal ring is given by:

$$B = \mu H$$

$$\phi = BA \quad \therefore B = \frac{\phi}{A}$$

$$B = \mu H$$

$$H = \frac{NI}{2\pi R}$$

$$\boxed{B = \frac{\mu NI}{2\pi R}}$$

$$\text{Total flux linkages} = N\phi$$

$$= NBA$$

$$= NA \frac{\mu NI}{2\pi R}$$

$$= \frac{\mu N^2 I A}{2\pi R}$$

Inductance of toroid is given by

$$L = \frac{\text{Total flux linkages}}{\text{Total current}}$$

$$L = \frac{\mu N^2 I A}{2\pi R I}$$

$$L = \frac{\mu N^2 A}{2\pi R} \text{ Henrys}$$

Where $A = \text{Area cross of toroid ring} = \pi r^2 \text{ (m}^2\text{)}$

for a toroid with No of turns (N) & $h = \text{height of the toroid with } (r_1) \text{ has inner radius and } (r_2) \text{ has outer radius, the inductance}$

- ce is given by $L = \frac{\mu N^2 h}{2\pi} \ln\left(\frac{r_2}{r_1}\right)$

$$L = \frac{\mu N^2 h}{2\pi} \ln\left(\frac{r_2}{r_1}\right)$$

Ex: calculate the inductance of a solenoid of 200 turns wound tightly on a cylindrical tube of 6cm diameter the length of the tube is 60cm & the solenoid is in air.

Given that

Turns $N = 200$

diameter $d = 6 \text{ cm} = 6 \times 10^{-2} \text{ m}$

length $l = 60 \text{ cm} = 60 \times 10^{-2} \text{ m}$

Solenoid inductance $L = \frac{\mu N^2 A}{l}$

$$L = \frac{\mu_0 \mu_r N^2 A}{l}$$

$$L = \frac{4\pi \times 10^{-7} \times (1) \times (200)^2 \times 9 \times 10^{-4}}{60 \times 10^{-2}}$$

$$A = \pi r^2$$

$$r = \frac{d}{2} = \frac{6 \times 10^{-2}}{2} = 3 \times 10^{-2}$$

$$r^2 = (3 \times 10^{-2})^2$$

$$= 9 \times 10^{-4}$$



$$L = \frac{4\pi \times 10^{-7} \times 4 \times 10^4 \times 9 \times 10^{-4}}{60 \times 10^{-2}}$$

$$L = \frac{\pi^2 \times 144 \times 10^{-7}}{6 \times 10^{-1}}$$

$$L = \frac{\pi^2 \times 144 \times 10^{-6}}{6}$$

$$L = \pi^2 \times 24 \times 10^{-6}$$

$$L = 2.368 \times 10^{-4} \text{ Henry}$$

pb: 2 A coil of 500 turns is wound on a closed iron ring of mean radius 10cm & cross sectional area of 3cm^2 find the Self inductance of the winding if the relative permibility of iron is 800

Give data

$$N = 500$$

$$\text{radius } R = 10\text{cm} = 10 \times 10^{-2}\text{m}$$

$$\text{Area } A = 3\text{cm}^2 = 3 \times (10^{-2})^2 = 3 \times 10^{-4}\text{m}^2$$

$$\mu_r \text{ relative permibility} = 800$$

$$L = \frac{\mu N^2 A}{2\pi R}$$

$$= \frac{\mu_0 \mu_r N^2 A}{2\pi R}$$

$$= \frac{4\pi \times 10^{-7} \times 800 \times 25 \times 10^4 \times 3 \times 10^{-4}}{2\pi \times 10^{-1}}$$

$$= 16 \times 25 \times 3 \times 10^{-4}$$

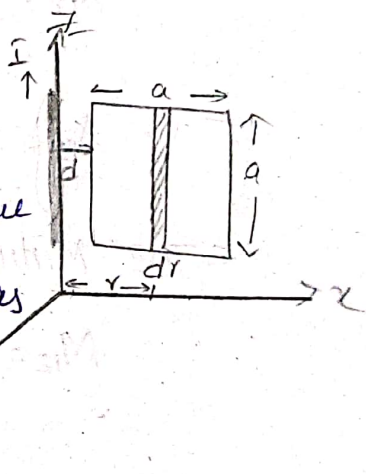
$$= 16 \times 75 \times 10^{-4}$$

Mutual inductance b/w a long straight wire & square loop lying in same plane

Consider a square loop with sides 'a', a straight long conductor is kept parallel to the longer side of the loop along z-axis.

Consider a long straight wire is circuit 1 while a square loop is circuit 2

The Magnetic field intensity at a distance of 'd' from long conductor due to current I_1 , it can be expressed as using Ampere circuit law.



$$\oint \vec{H} \cdot d\vec{L} = I$$

$$H_1(2\pi r) = I_1$$

$$H_1 = \frac{I_1}{2\pi r} \hat{\phi} \rightarrow (1)$$

Magnetic flux density

$$B_1 = \mu H_1$$

$$B_1 = \mu_0 \mu_r H_1$$

$$B_1 = \mu_0 \mu_r \frac{I_1}{2\pi r} \hat{\phi} \rightarrow (2)$$

The flux linkages in circuit 2 due to current in circuit -1

$$\lambda_{12} = \int_{S_2} B_1 \cdot ds_2$$

Since $ds_2 = a \cdot dr \cdot a$

$$\lambda_{12} = \int_d^{d+a} \frac{\mu I_1}{2\pi r} a \psi \cdot (adr \cdot d\psi)$$

$$\lambda_{12} = \frac{\mu I_1 a}{2\pi} \int_d^{d+a} \frac{1}{r} dr$$

$$\lambda_{12} = \frac{\mu I_1 a}{2\pi} \log_e \frac{d+a}{d}$$

$$\lambda_{12} = \frac{\mu I_1 a}{2\pi} (\log_e(d+a) - \log_e d)$$

$$\lambda_{12} = \frac{\mu I_1 a}{2\pi} \log_e \left(\frac{d+a}{d} \right)$$

$$\lambda_{12} = \frac{\mu I_1 a}{2\pi} \log_e \left(1 + \frac{a}{d} \right)$$

∴ The Mutual inductance

$$M_{12} = \frac{\lambda_{12}}{I_1} = \frac{\mu I_1 a}{2\pi} \log_e \left(1 + \frac{a}{d} \right) \cdot \frac{1}{I_1}$$

$$= \frac{\mu a}{2\pi} \ln \left(1 + \frac{a}{d} \right)$$

(or)

Mutual inductance b/w a long wire straight wire & rectangular loop in same plane $M_{12} =$

$$M_{12} = \frac{\lambda_{12}}{I_1} = \frac{\mu b}{2\pi} \ln \left(1 + \frac{a}{d} \right)$$

3 Energy stored in a magnetic field

Consider a differential volume in magnetic field B , consider at the top & bottom surfaces of a differential volume conducting sheets ΔI are present since energy stored magnetic field

$$E = \frac{1}{2} L I^2$$

The inductance of the conductor

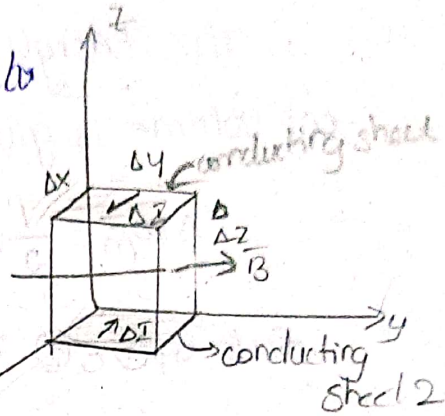
$$L = \frac{N\Phi}{I}$$

$$\therefore N = 1$$

$$L = \frac{\Phi}{I}$$

The differential inductance

$$\Delta L = \frac{\Delta\Phi}{\Delta I}$$



The Magnetic flux

$$\Phi = BA$$

$$\Delta\Phi = B\Delta S$$

$$\Delta L = \frac{B\Delta S}{\Delta I}$$

where $\Delta S =$ differential surface area $= \Delta x \Delta z$

$$\therefore \Delta L = \frac{B\Delta x \Delta z}{\Delta I}$$

$$B = \mu H$$

$$\Delta L = \frac{\mu H \Delta x \Delta z}{\Delta I} \rightarrow \textcircled{1}$$

The differential current ΔI can be in terms of magnetic field intensity H . The current flowing through the conduction sheet present at the top & bottom is in y -direction.

$$\int H \cdot dL = I$$

$$H \cdot dL = \Delta I$$

$$\text{Since } \Delta L = \Delta y$$

$$\therefore \Delta I = H \Delta y \rightarrow \textcircled{2}$$

∴ The energy stored in a inductance of a different volume is given by $\Delta W_m = \frac{1}{2} \Delta L (\Delta I)^2$

$$\Delta W_m = \frac{1}{2} \Delta L (\Delta I)^2 \rightarrow (3)$$

Sub eq (1) & (2) in eq (3) we get

$$\Delta W_m = \frac{1}{2} \frac{\mu H \Delta x \Delta z}{\Delta l} (H \Delta y)^2$$

$$= \frac{1}{2} \frac{\mu H \Delta x \Delta z}{(H \Delta y)} (H \Delta y)^2$$

$$= \frac{1}{2} (\mu H^2 \Delta x \Delta y \Delta z)$$

$$= \frac{1}{2} \mu H^2 \Delta x \Delta y \Delta z$$

$$\therefore \Delta V = \Delta x \Delta y \Delta z$$

$$\Delta W_m = \frac{1}{2} \mu H^2 \Delta V \rightarrow (4)$$

The magneto statics Energy density is defined as

$$w_m = \lim_{\Delta V \rightarrow 0} \frac{\Delta W_m}{\Delta V}$$

$$w_m = \lim_{\Delta V \rightarrow 0} \frac{\frac{1}{2} \mu H^2 \Delta V}{\Delta V}$$

$$w_m = \frac{1}{2} \mu H^2 \text{ (Joules/m}^3\text{)}$$

The magneto statics Energy density can be expressed in different forms

$$w_m = \frac{1}{2} (\mu H) H$$

$$w_m = \frac{1}{2} BH$$

$$W_m = \frac{1}{2} B \left(\frac{B}{\mu} \right)$$

$$W_m = \frac{1}{2} \frac{B^2}{\mu}$$

$$\boxed{W_m = \frac{1}{2} \left(\frac{B^2}{\mu} \right)}$$

$$B = \mu H$$

$$H = \frac{B}{\mu}$$

5. Time Varying field

Faradays laws of Electro Magnetic induction:-

First law:-

Whenever a conductor cuts the magnetic field or magnetic flux lines then dynamically induced emf in that conductor

Second law:-

Whenever the dynamically induced emf is directly proportional to the rate of change of flux linkages

$$e \propto \frac{d\phi}{dt} \text{ volts}$$

$$e = -N \frac{d\phi}{dt}$$

Explanation:-

Let us consider a coil having a no. of turns 'N', and flux linkages with coil minimum ϕ_1 & maximum ϕ_2 and the conductor in a time period 't' in seconds

The initial flux linkages = $N\phi_1$

The final flux linkage = $N\phi_2$

change of flux linkages = $N\phi_2 - N\phi_1$

$$\phi = N(\phi_2 - \phi_1)$$

Rate of change of flux linkages = $N \frac{d\phi}{dt}$

$$\therefore e = -N \frac{d\phi}{dt}$$

The Negative is due to lenz's law & indicates

The induced voltage is direction to the opposes to change in flux linkage that producing it

2. Poynting Vector & Poynting theorem:

*] In Electromagnetic waves, an energy can be transported from the transmitter to the Receiver.

*] The energy stored in an electric field and magnetic field energy is transmitted at a certain rate of energy flow which can be calculated with the help of Poynting theorem.

*] The product $\vec{E} \times \vec{H}$ gives the new quantity which is expressed as watt/unit area. Thus this quantity is called power density.

*] Consider that the field is transmitted in the form of an electromagnetic waves from an antenna. Both the fields are sinusoidal in nature.

*] The power radiated from the antenna as a particular direction so the power density is given by

$$\vec{P} = \vec{E} \times \vec{H}$$

where \vec{P} is called Poynting vector.

\vec{E} is called electric field.

*] Poynting theorem is based on law of conservation of energy in Electro Magnetism.

*] The net power flowing out of a given volume (V) is equal to the rate of decrease in the

energy stored within volume V - The omic power dissipated

*] Since $\vec{E} = E_x \cdot \vec{a}_x$

*] $\vec{H} = H_y \cdot \vec{a}_y$

*] $\vec{P} = \vec{E} \times \vec{H}$

$\vec{P} = E_x \cdot \vec{a}_x \times H_y \cdot \vec{a}_y$

$\therefore \vec{a}_x \times \vec{a}_y = \vec{a}_z$

$\vec{P} = E_x H_y \vec{a}_z$

$\vec{P} = P_z \vec{a}_z$

$\therefore P_z = E_x H_y$

*] The above equation \vec{E} , \vec{P} , \vec{H} are mutually perpendicular to each other.

*] Consider that the electric field propagates in free space is given by $\vec{E} = E_m (\cos(\omega t - \beta z)) \vec{a}_x$

*] The ratio of magnitudes of \vec{E} & \vec{H} depends on its intrinsic impedance η for free space

$\eta = \eta_0 = \frac{E_m}{H} = 120\pi = 377 \Omega$

*] The electromagnetic wave travels at speed of light in

$\vec{H} = H_m (\cos(\omega t - \beta z)) \vec{a}_y$

*] Since $H_m = \frac{E_m}{\eta_0}$

$\vec{H} = \frac{E_m}{\eta_0} (\cos(\omega t - \beta z)) \vec{a}_y$

$\vec{P} = \vec{E} \times \vec{H}$

$\vec{P} = E_m (\cos(\omega t - \beta z)) \vec{a}_x \times \frac{E_m}{\eta_0} (\cos(\omega t - \beta z)) \vec{a}_y$

$\vec{P} = \frac{E_m^2}{\eta_0} \cos^2(\omega t - \beta z) \vec{a}_z \text{ (watt/m}^2\text{)}$

3 Maxwell equations for Time Variant fields:-

(a) Maxwell Eqn from Faraday's Law:

consider Faraday's which relates EMF induced in a circuit to the time rate of decrease of total Magnetic flux linking the circuit.

$$\oint \vec{E} \cdot d\vec{l} = - \int \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s}$$

*] This equation is called as Maxwell equation derived from Faraday's law expressed in integral form

*] The Total Electromotive force induced in a closed path is equal to the negative surface integral of the rate of change of flux density w.r.t time over an entire surface bounded by the same closed path.

$$\int_S (\nabla \times \vec{E}) \cdot d\vec{s} = - \int_S \frac{\partial \vec{B}}{\partial t} \cdot d\vec{s}$$

$$\nabla \times \vec{E} = - \frac{\partial \vec{B}}{\partial t}$$

This eqn is called Maxwell equation derived from Faraday's expressed in point form differential form

(b) Maxwell equation derived from Ampere circuit law:-

According to Ampere circuit law the line integral Magnetic field intensity \vec{H} around a closed path is equal to the current

enclosed by $\oint \vec{H} \cdot d\vec{l} = \vec{i}$ enclosed
Statement:

The Total magnetomotive force around any closed path is equal to the surface integral of the conduction and displacement current densities over the entire surface bounded by the same closed path

$$\oint \vec{H} \cdot d\vec{l} = \int \vec{J} \cdot d\vec{s}$$

$$\oint \vec{H} \cdot d\vec{l} = \int \left[\vec{J} + \frac{\partial \vec{D}}{\partial t} \right] \cdot d\vec{s}$$

This equation is called Maxwell Equation derived from Ampere circuit Law in integral form

$$\int (\nabla \times \vec{H}) \cdot d\vec{s} = \int \left[\vec{J} + \frac{\partial \vec{D}}{\partial t} \right] \cdot d\vec{s}$$

$$\nabla \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$$

The equation is called as Maxwell equation in point form or differential form derived from Amperes ckt law